

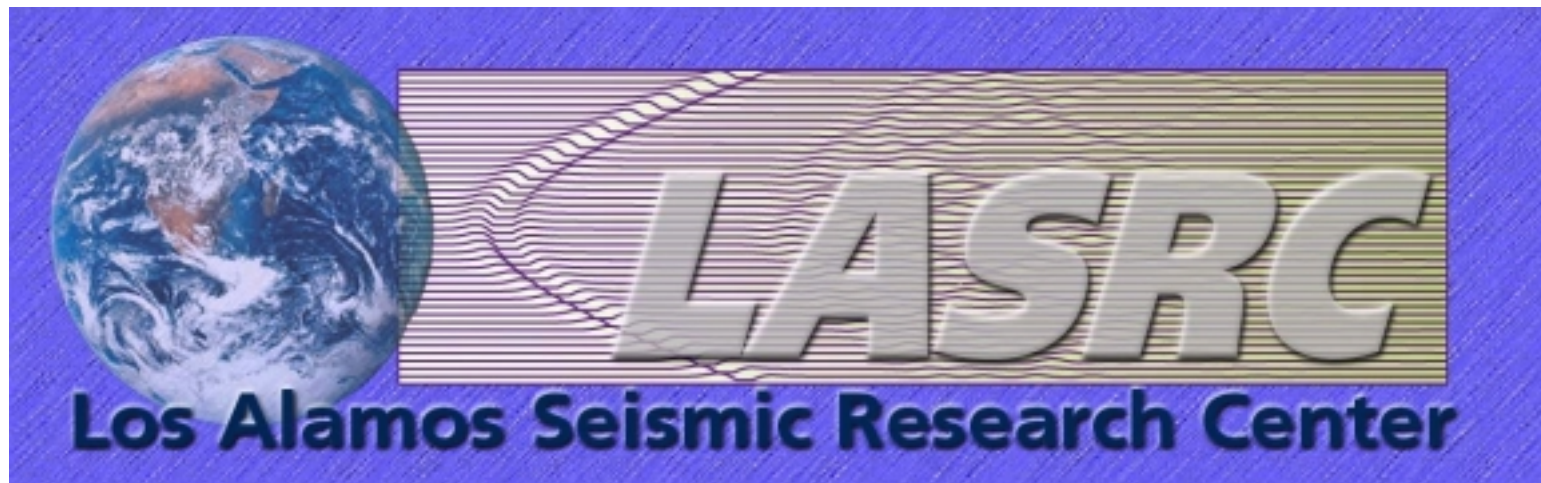
Numerical Study of Envelope Broadening in Random Media

by

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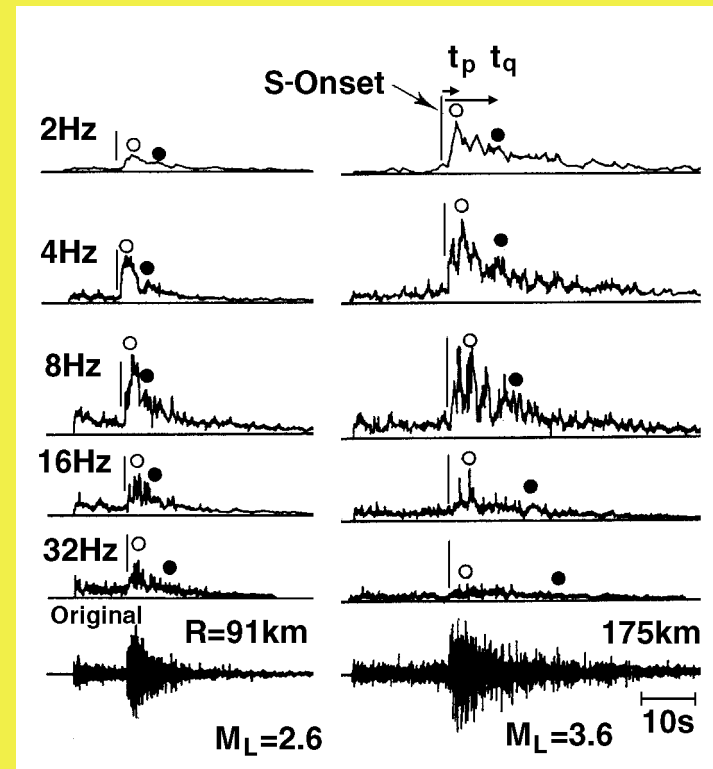
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Forward Propagation and Envelope Broadening

Regional seismograms of deep earthquakes in Japan have long coda duration due to multiple forward scattering

Lg waveforms influenced by multiple forward scattering



Seismograms from Deep Earthquakes in Japan

Model for Multiple Forward Scattering

Forward scattering dominates when characteristic length of variation in medium velocity (or density) is $>$ wavelength of waves

Can model with *Parabolic Wave Equation* which neglects the second derivative in the wave equation with respect to the global ray direction

Split-Step Fourier Modeling (Stoffa, et al. 1990)

Extended Local Rytov Approximation (Huang et al, 1999)

Markov Approximation



Objectives

Investigate envelope broadening using numerical (Rytov and Finite Difference) modeling

Compare results obtained using Markov Approximation with numerical results obtained using finite difference and approximate method based on Rytov Approximation.

Investigate variations in waveforms obtained from various realizations of random media

Investigate differences obtained using various modeling approaches



Previous Numerical Modeling of Wave Propagation in Random Media

Frankel and Clayton (1986, *JGR*)

focused on near-source region (backscattering)

Ikelle, Young and Daube (1993, *Geophysics*)

Focused on effects of anisotropy in random media
characterization



Random Media Models

- Choose Model Where Forward Scattering Dominates
- Propagation to Regional Distances (>100 km)
- Write velocity as $V(\mathbf{x}) \equiv V_0 + \delta V(\mathbf{x}) = V_0(1 + \xi(\mathbf{x}))$
 $\xi(\mathbf{x})$ is fractional fluctuation of wave velocity
 V_0 is chosen so that $V_0 = \langle V(\mathbf{x}) \rangle$ and $\langle \xi(\mathbf{x}) \rangle = 0$



Random Media Models

Autocorrelation function (ACF) of the medium

$$R(\mathbf{x}) \equiv \langle \xi(\mathbf{y}) \xi(\mathbf{y} + \mathbf{x}) \rangle$$

Magnitude of fractional fluctuation is mean square (MS) fractional fluctuation:

$$\varepsilon^2 \equiv R(0) = \langle \xi(\mathbf{x})^2 \rangle$$

Choose Gaussian ACF

$$R(\mathbf{x}) = R(r) = \varepsilon^2 e^{-\frac{x^2}{a_x^2}} e^{-\frac{z^2}{a_z^2}}$$

where a_x and a_z are correlation distances.

Medium characterized by

RMS fractional fluctuation ε

Correlation distance a

We choose $\varepsilon = .05$

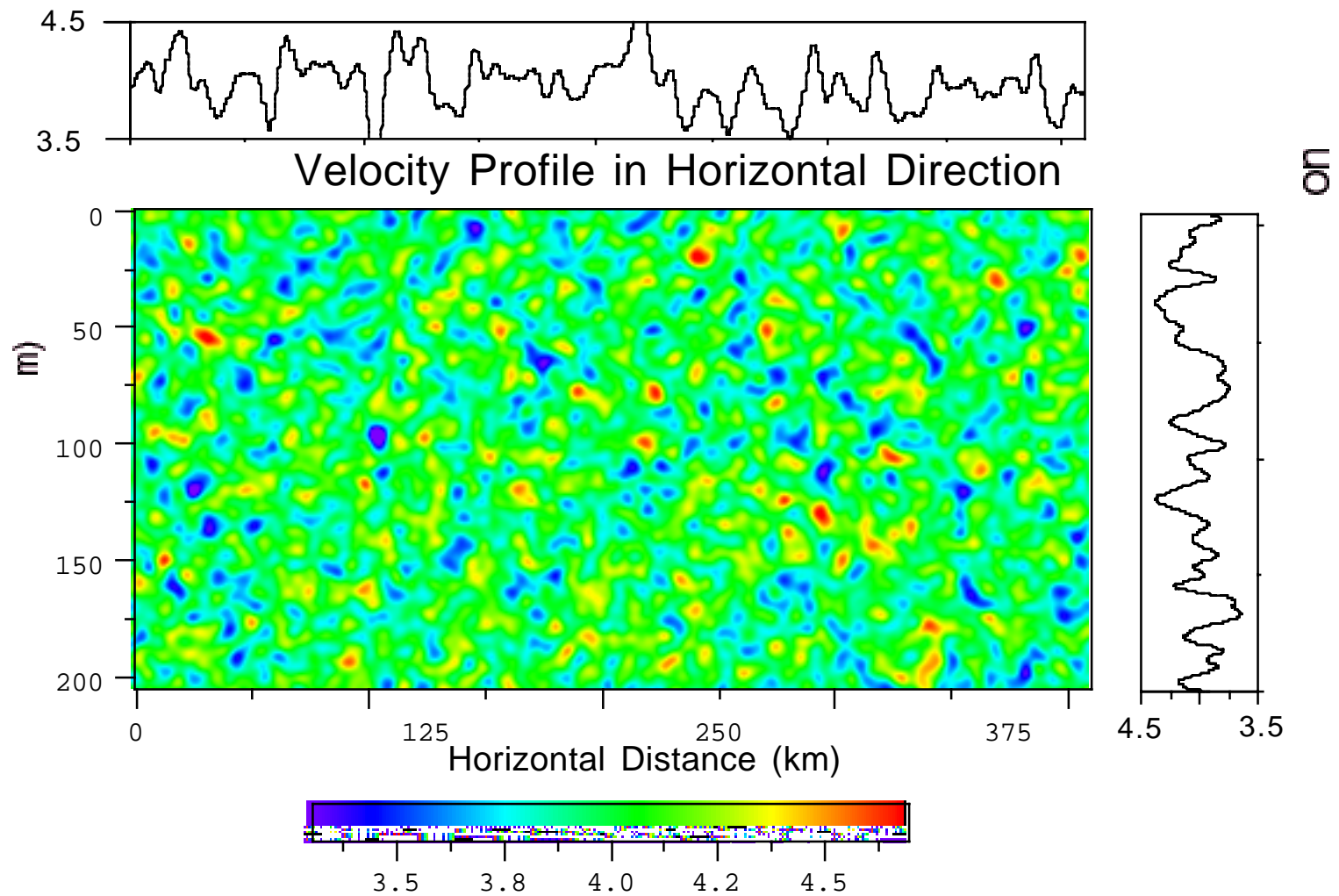
$$a = 5 \text{ km}$$

$$V_0 = 4. \text{ km/s}$$

Source: Ricker time history with 2Hz
dominant frequency

Wavelength in background medium is 2 km

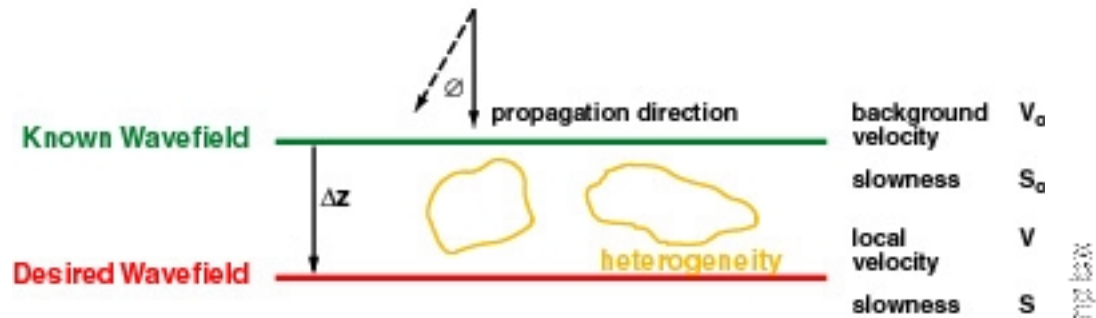
Gaussian Random Medium



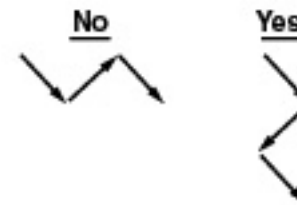
Markov Approximation

- Method to calculate mean wavefield for propagation in random media when forward scattering dominates
- Lee and Jakopi (1975) Sreenivasiah et al. (1976) used to calculate envelope for plane waves incident on 3D random media
- Sato (*JGR*, 1989) used to study envelope broadening
- Here, calculated for point source in 2D random media

One-Way Wave Propagation



Reverberations not modeled
Modeling is fast and local (less memory)



Our methods use
Local Born
Local Rytov

$$P = P_0 + P_S$$

$$P = P_0 \rho^{0s} \quad \rho_s \text{ complex}$$

Approximation

⇒ Background Velocity V_0 can be unique for each layer

→ *Background Velocity V_0 can be unique for each layer*

Local Rytov Fourier Method

(Huang et al. ,*Geophysics* **64**, 1535-1545, 1999)

Wave Equation

$$\left(\Delta - \frac{1}{V(\mathbf{x})^2} \partial_t^2 \right) u(\mathbf{x}, t) = 0$$

$$\Delta = \partial^2 / \partial x^2 + \partial^2 / \partial z^2$$

Using

$$V(\mathbf{x}) \equiv V_0 + \delta V(\mathbf{x}) = V_0 (1 + \xi(\mathbf{x}))$$

$$|\xi| \ll 1$$

We get

$$\left(\Delta - \frac{1}{V_0^2} \partial_t^2 \right) u(\mathbf{x}, t) + \frac{2}{V_0^2} \xi(\mathbf{x}) \partial_t^2 u(\mathbf{x}, t) = 0$$

Frequency Domain

$$\left(\Delta + \frac{\omega^2}{V_0^2} \right) u(\mathbf{x}, \omega) = 2 \frac{\omega^2}{V_0^2} \xi(\mathbf{x}) u(\mathbf{x}, \omega)$$

Local Rytov Fourier Method (cont.)

Solution of form

$$u(\mathbf{x}, \omega) = e^{\phi_0(\mathbf{x}, \omega) + \phi_s(\mathbf{x}, \omega)}$$

Homogeneous
propagation

$$u_0(\mathbf{x}, \omega) = e^{\phi_0(\mathbf{x}, \omega)}$$

Extrapolate wavefield
in primary propagation
direction

$$u(x, z_{i+1}; \omega) = u_0(x, z_{i+1}; \omega) e^{\phi_s(x, z_{i+1}; \omega)}$$

$$u_0(x, z_{i+1}; \omega) = F_{k_x}^{-1} \left\{ e^{ik_z \Delta z} F_x \{u(x, z_i, \omega)\} \right\}$$

Local Rytov Fourier Method (cont.)

Scattering term: $\phi_S(x, z_{i+1}; \omega) \equiv \psi(x, z_{i+1}; \omega) / u_0(\mathbf{x}_T, z_{i+1}; \omega)$

$$\psi(x, z_{i+1}; \omega) = F_{k_x}^{-1} \left\{ \frac{k_0}{k_z} e^{ik_z \Delta z} F_x \{ i\omega \Delta z \Delta s(x, z_i) u(x, z_i; \omega) \} \right\}$$

Slowness perturbation: $\Delta s(x, z_i) = -\xi(x, z_i) / V(x, z_i)$

Wavenumbers: $k_z = \sqrt{k_0^2 - k_x^2}$, $k_0 = \omega / V_0$

Fourier Transforms: $F_x, F_{k_x}^{-1}$

Finite Difference Modeling

- 2D Finite Difference Code
- 4th order in Space
- 2th Order in Time
- Holberg Coefficients to Minimize Dispersion
- Absorbing Boundaries

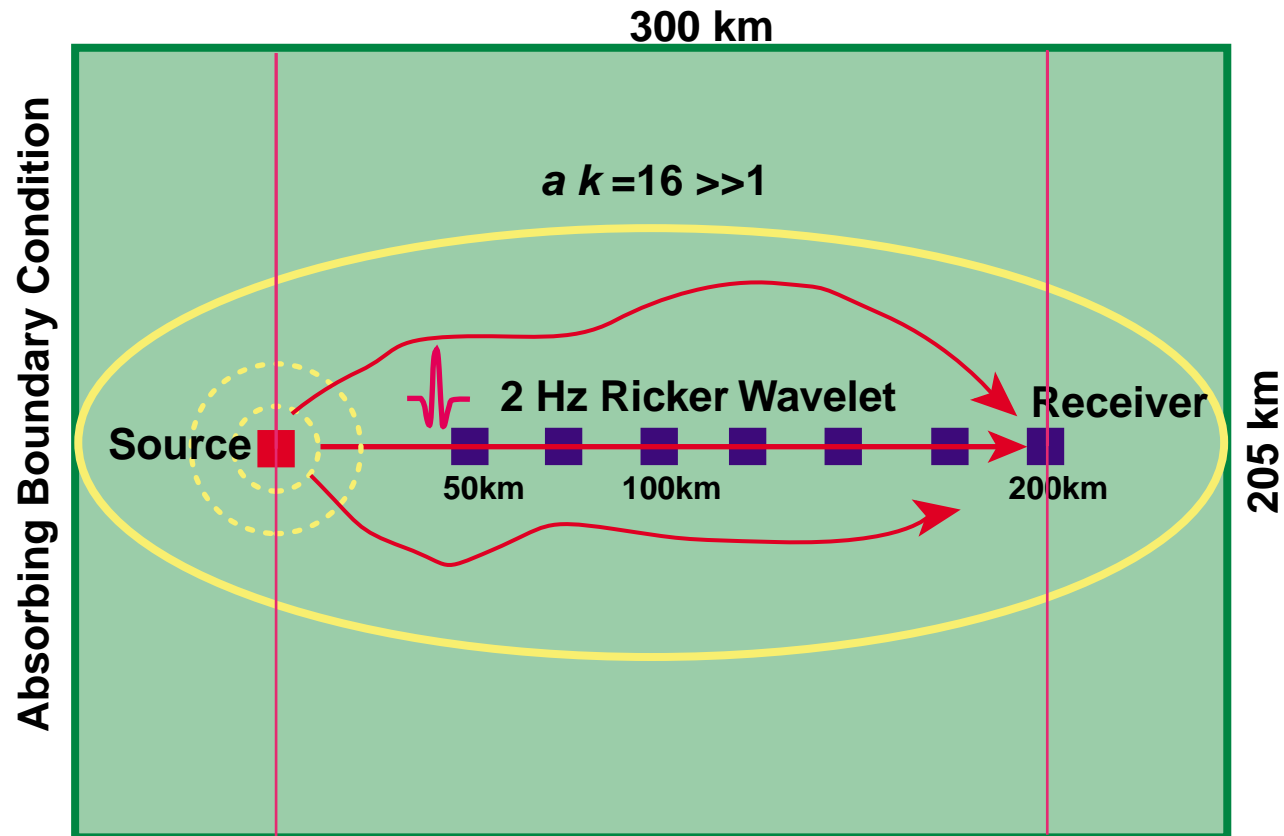


Numerical Experiment

Finite Difference: Grid Size 50m, Time Step 4 ms

Random Media of Gaussian ACF

$\varepsilon = 0.05$, $a = 5$ km, $V_0 = 4$ km/s



4,096 x 6,000 grid points in region of validity

Finite Difference Grid: 4,146 x 6,052

Extended Local Rytov Fourier Grid: 4,096 X 4,096

Valid up to 25s after the onset at 200 km distance

Finite Difference Seismograms in Homogeneous Media

Velocity = 4 km/s; Source: 2 Hz Ricker

Trace Length = 60 s; Trace Amplitudes Normalized

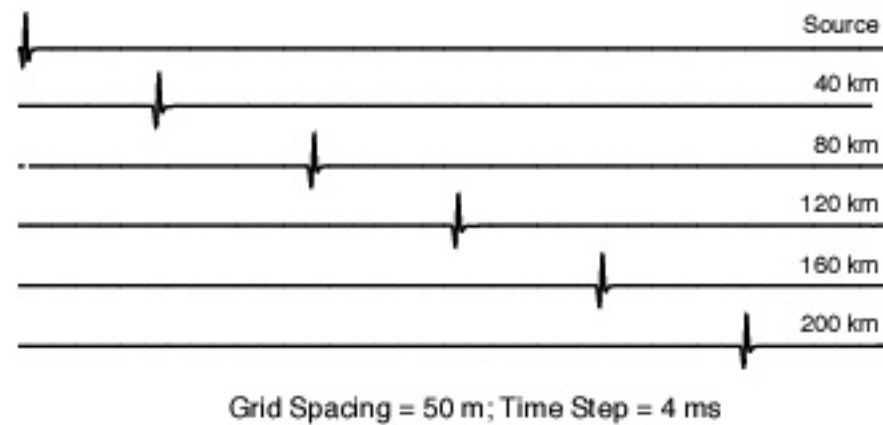
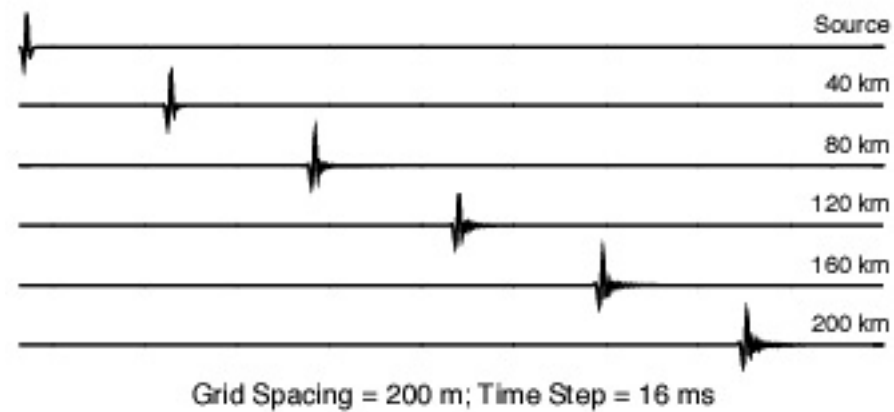
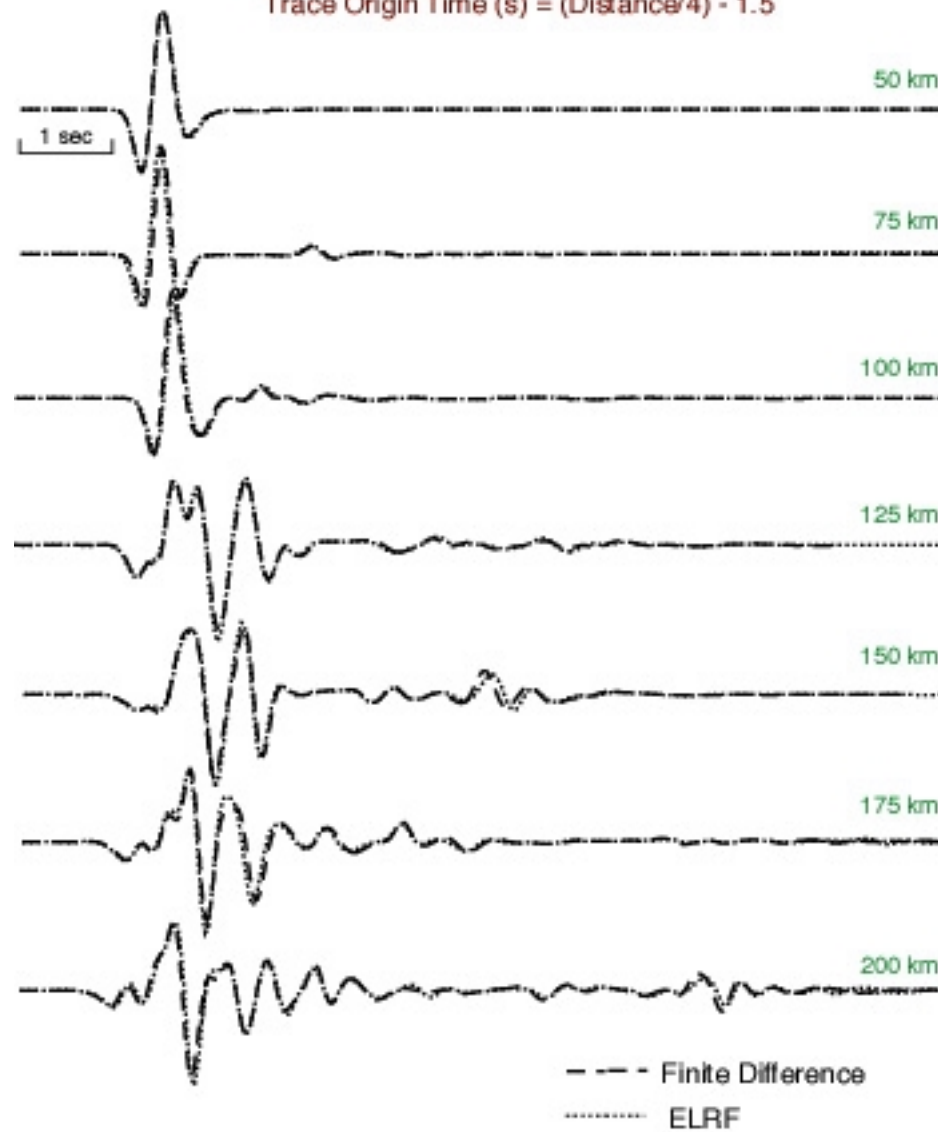


Figure 2

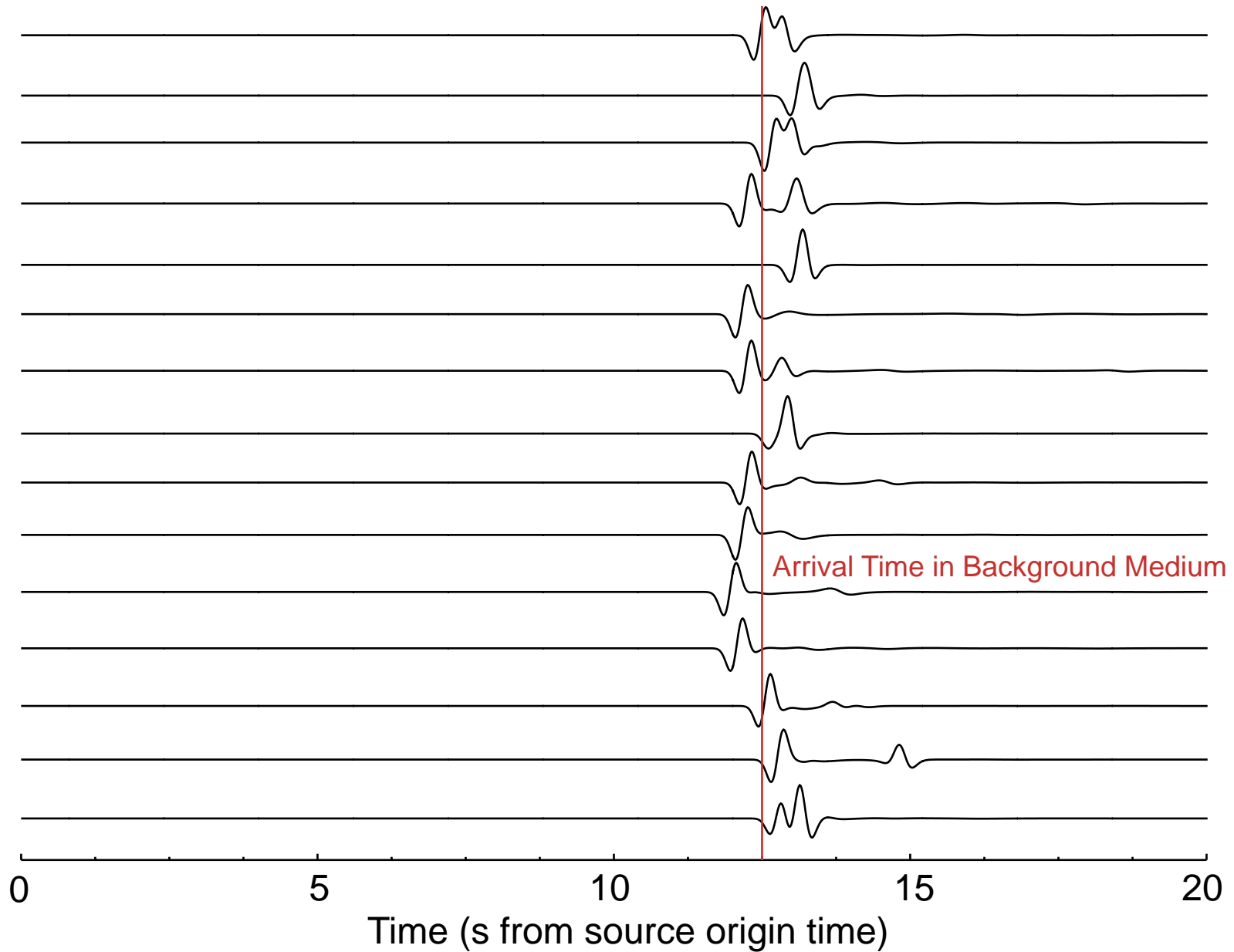
Comparison of Traces Calculated Using Finite Difference and ELRF

10 s shown for each trace

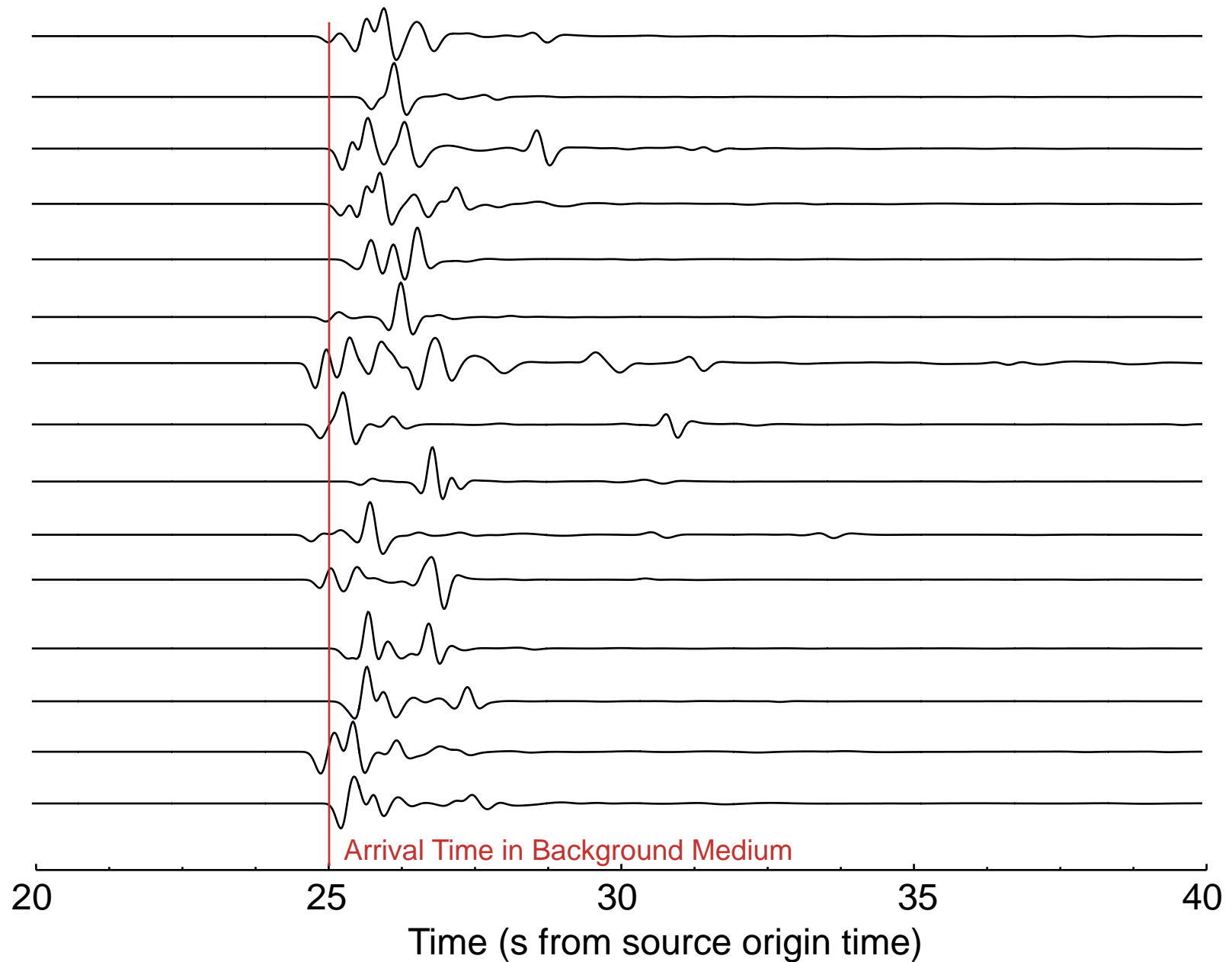
$$\text{Trace Origin Time (s)} = (\text{Distance}/4) - 1.5$$



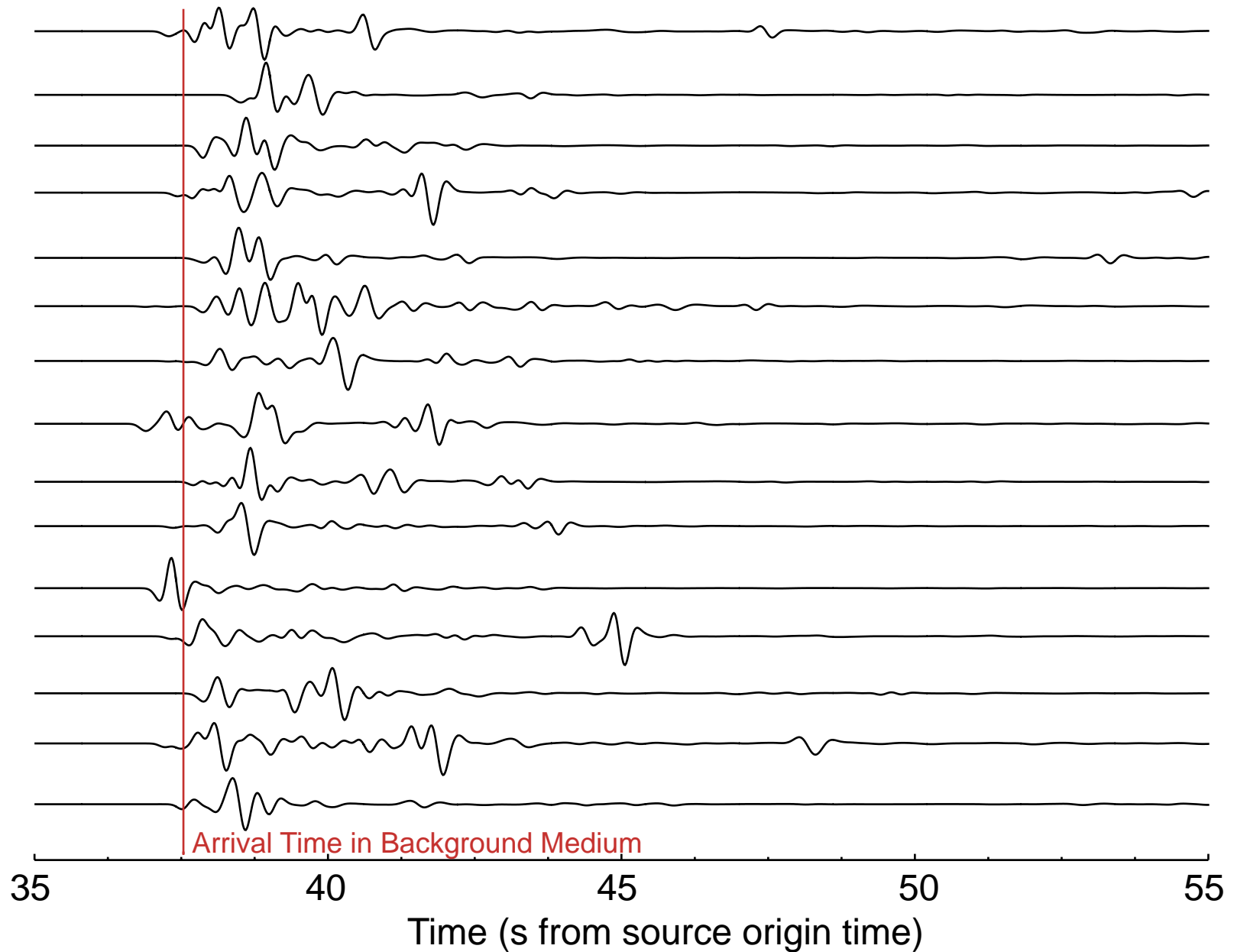
Waveforms at 50 km from Source in 15 Realizations of Large Model



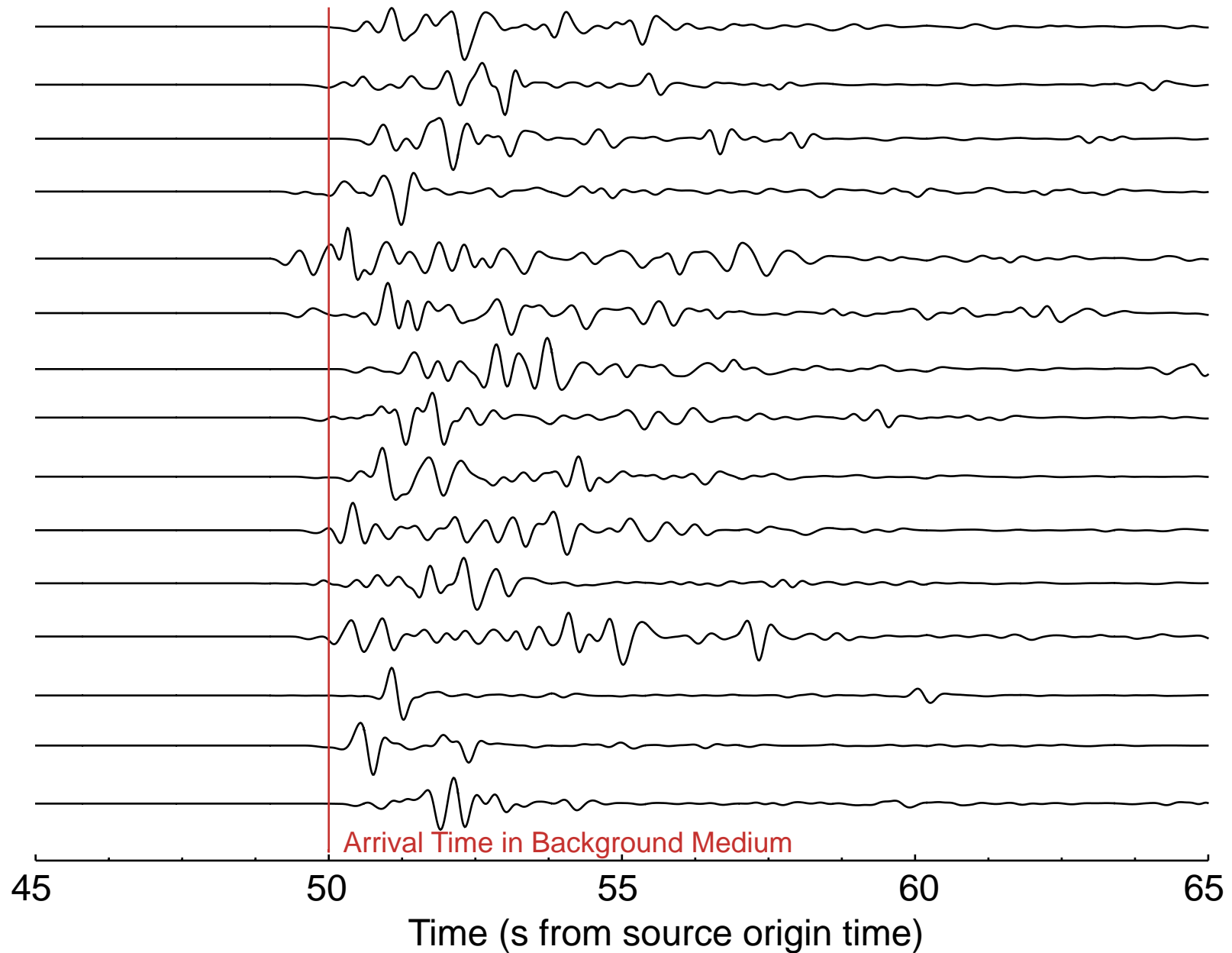
Waveforms at 100 km from Source in 15 Realizations of Large Model



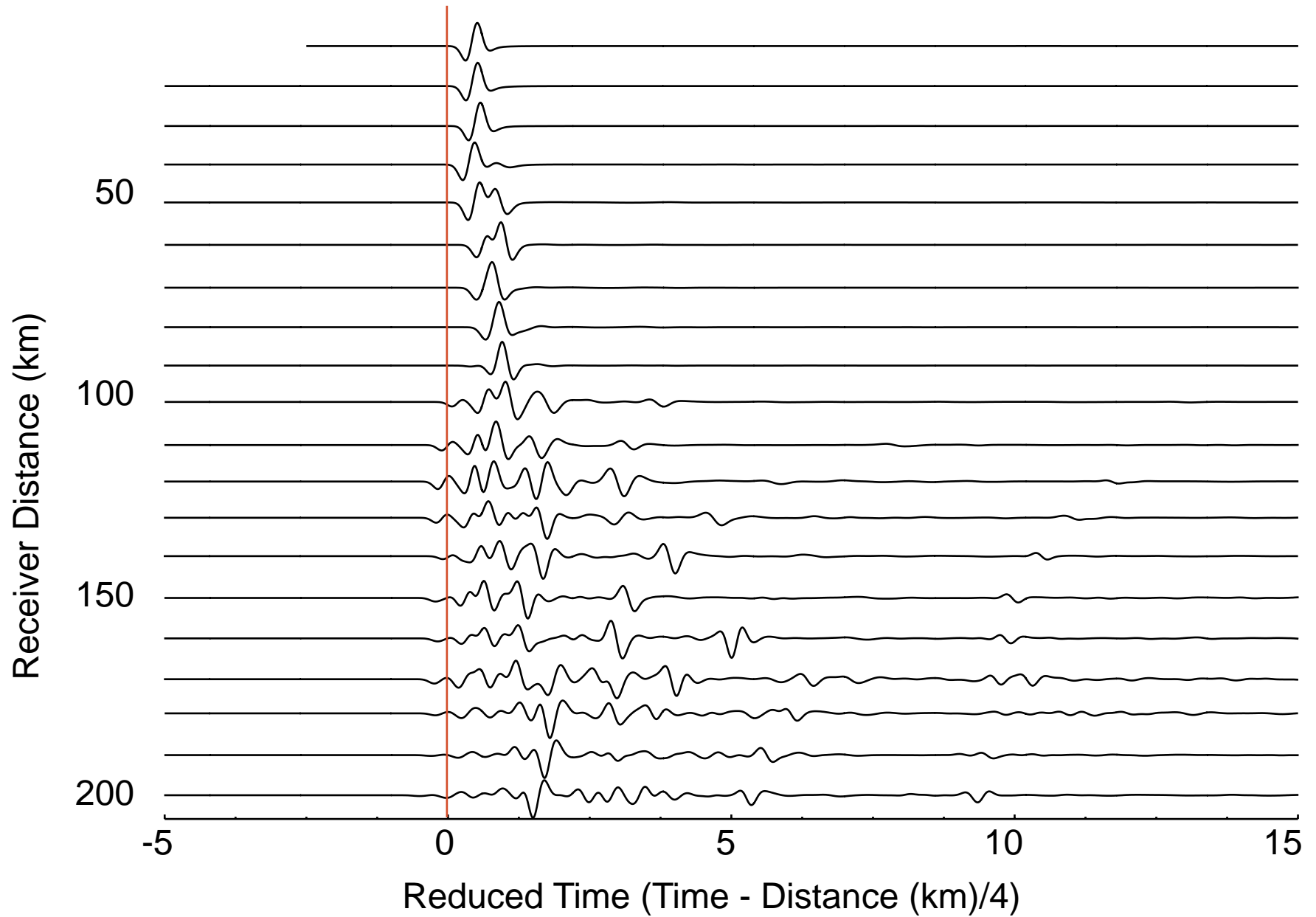
Waveforms at 150 km from Source in 15 Realizations of Large Model



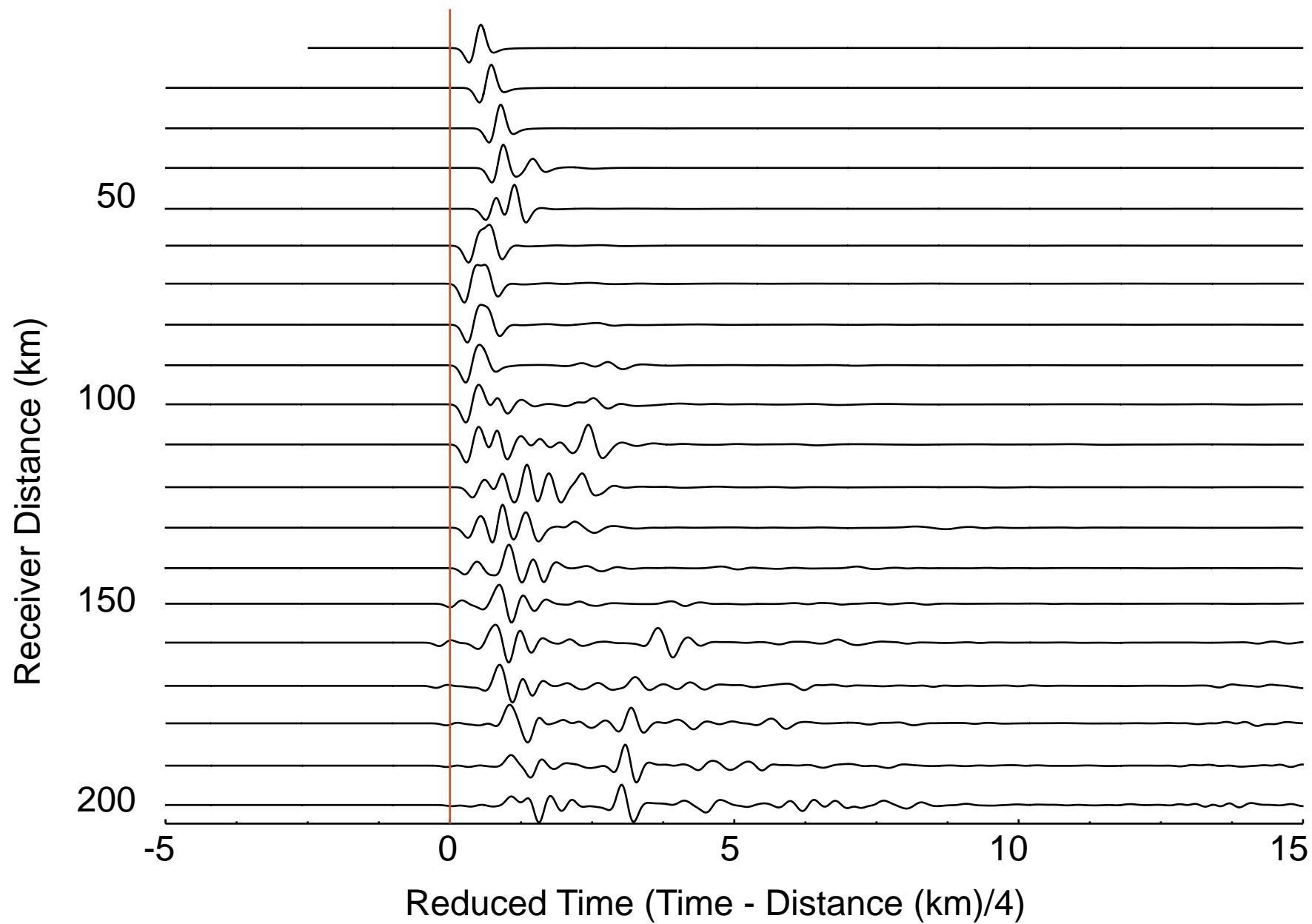
Waveforms at 200 km from Source in 15 Realizations of Large Model



Record Section from Finite Difference of Large Model (15500)
Trace Amplitudes Scaled Independently



Record Section from Finite Difference of Large Model (16900)
Trace Amplitudes Scaled Independently



Ensemble Average

- Calculate Waveforms for Propagation in many realizations of media
- Find mean envelope Shape
 - Sum Square Traces from each realization
 - Smooth over .32s
 - Plot square root of result

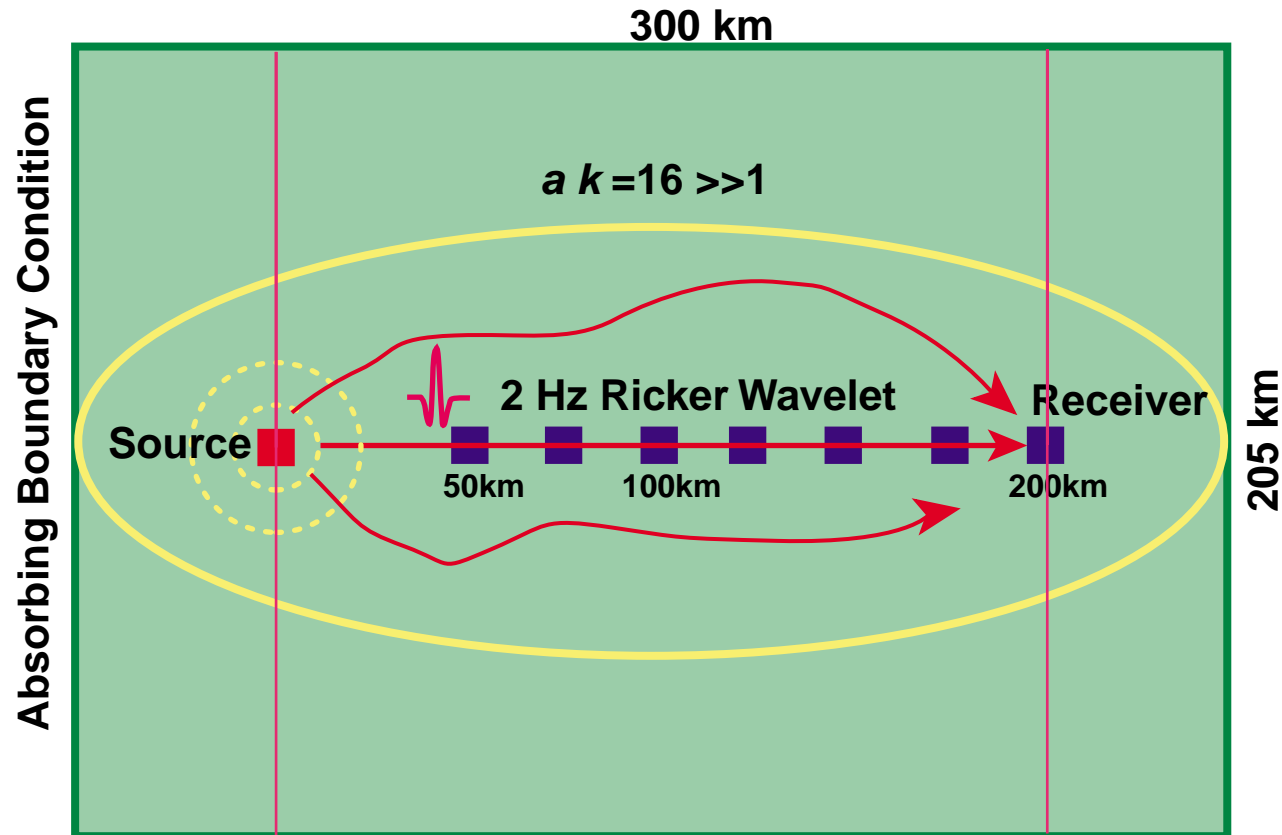


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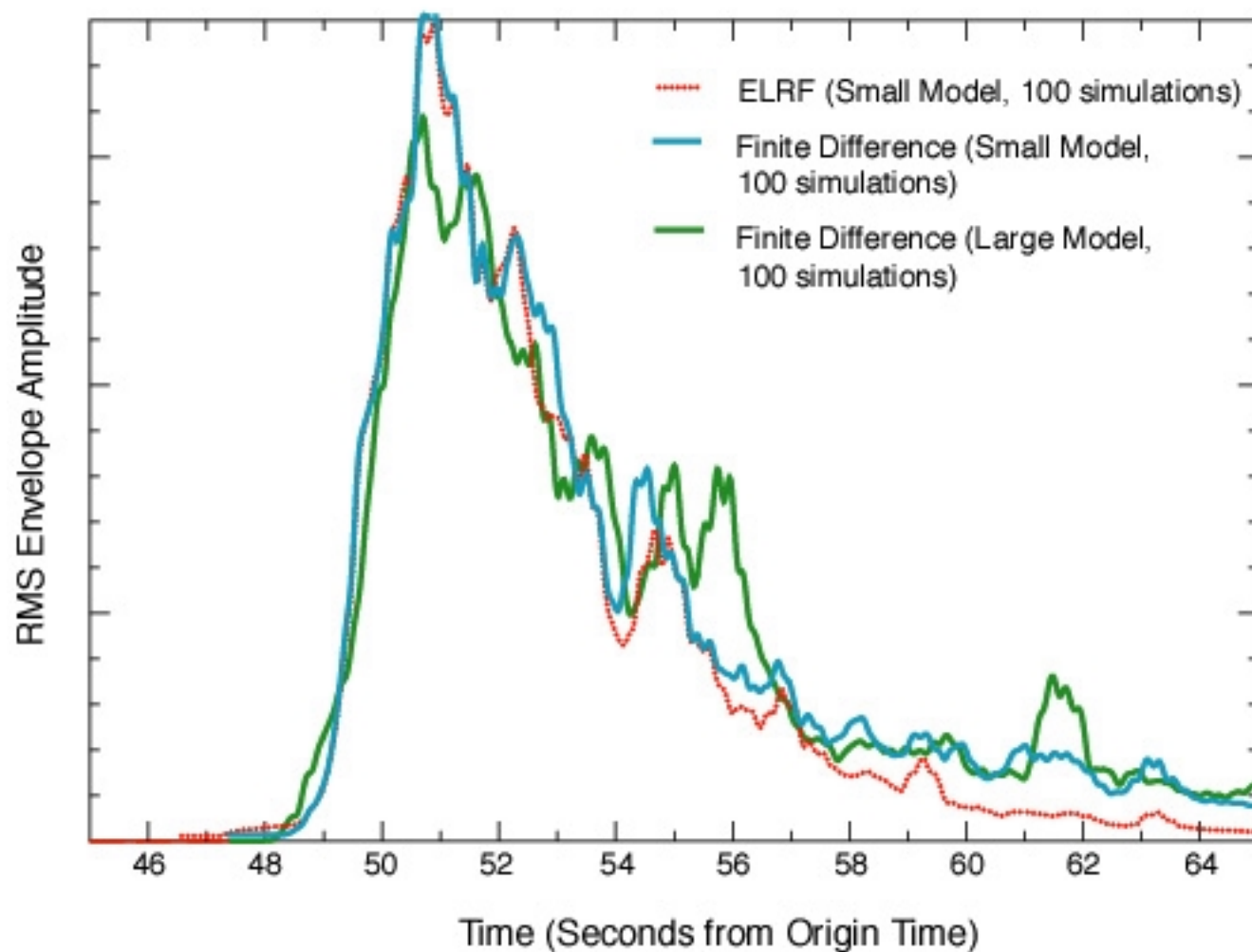
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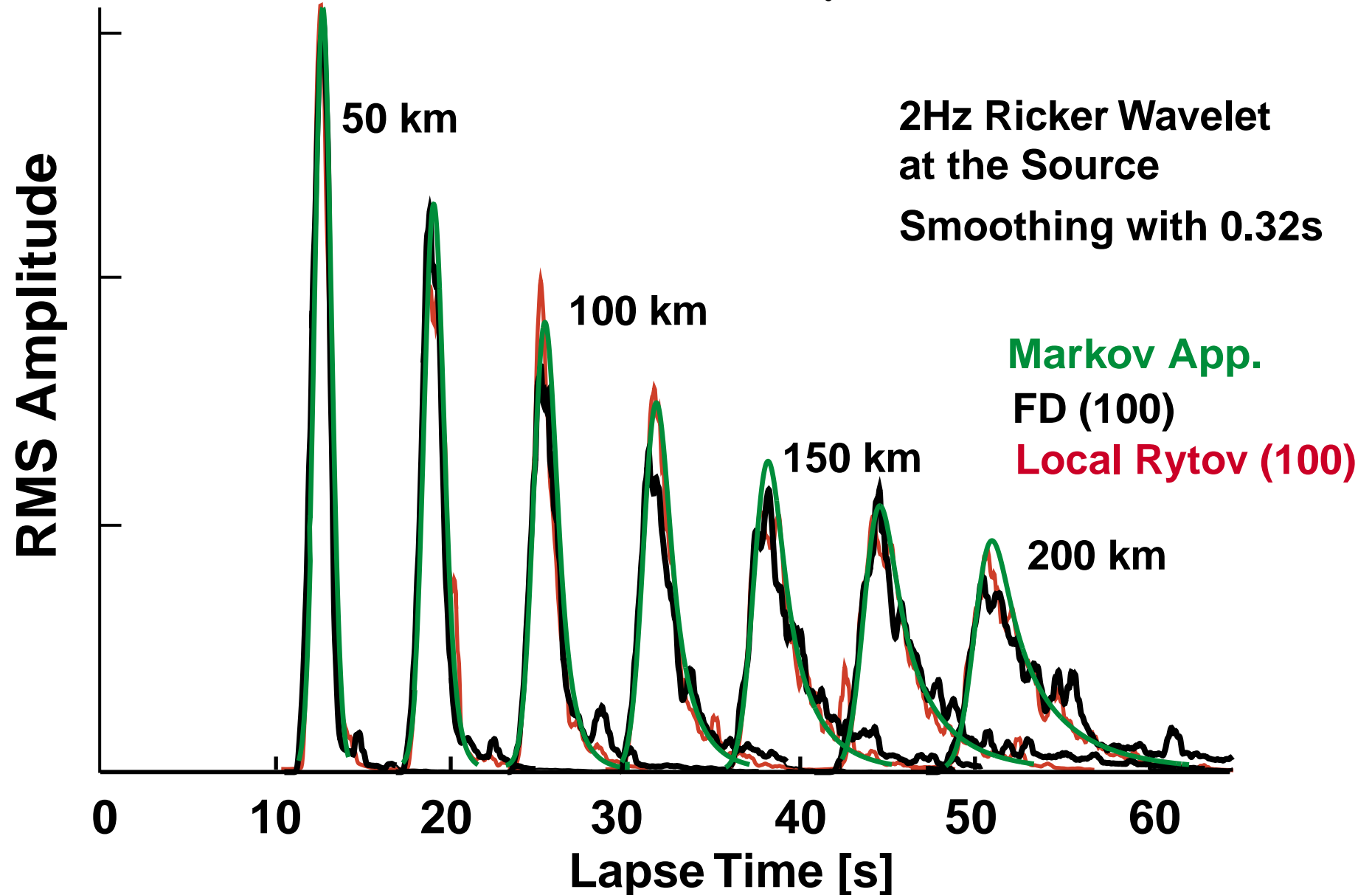
Valid up to 25s after the onset at 200 km distance

RMS Envelopes at 200 km Calculated Using Different Methods

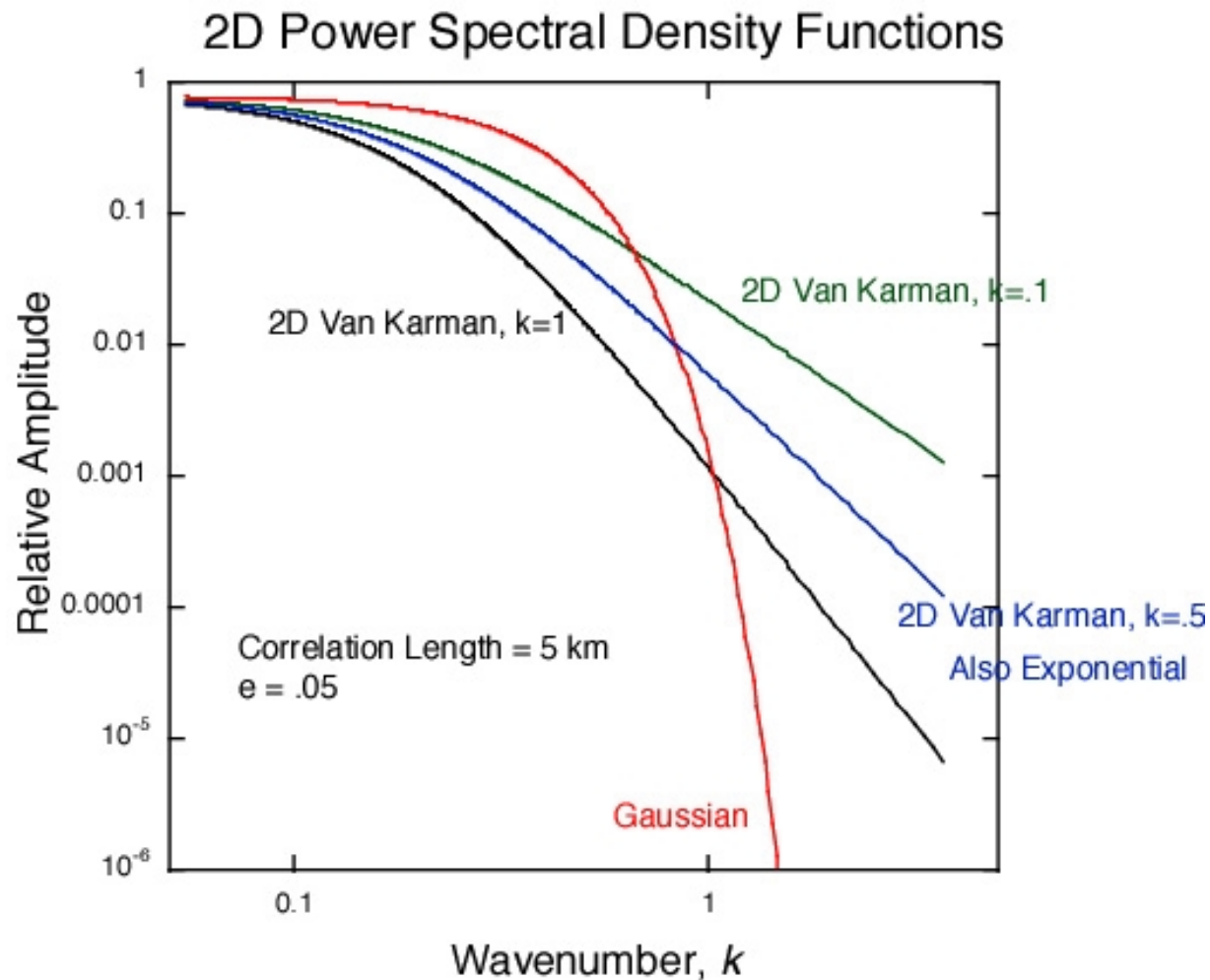


Comparison of RMS Envelopes in 2-D Random Media

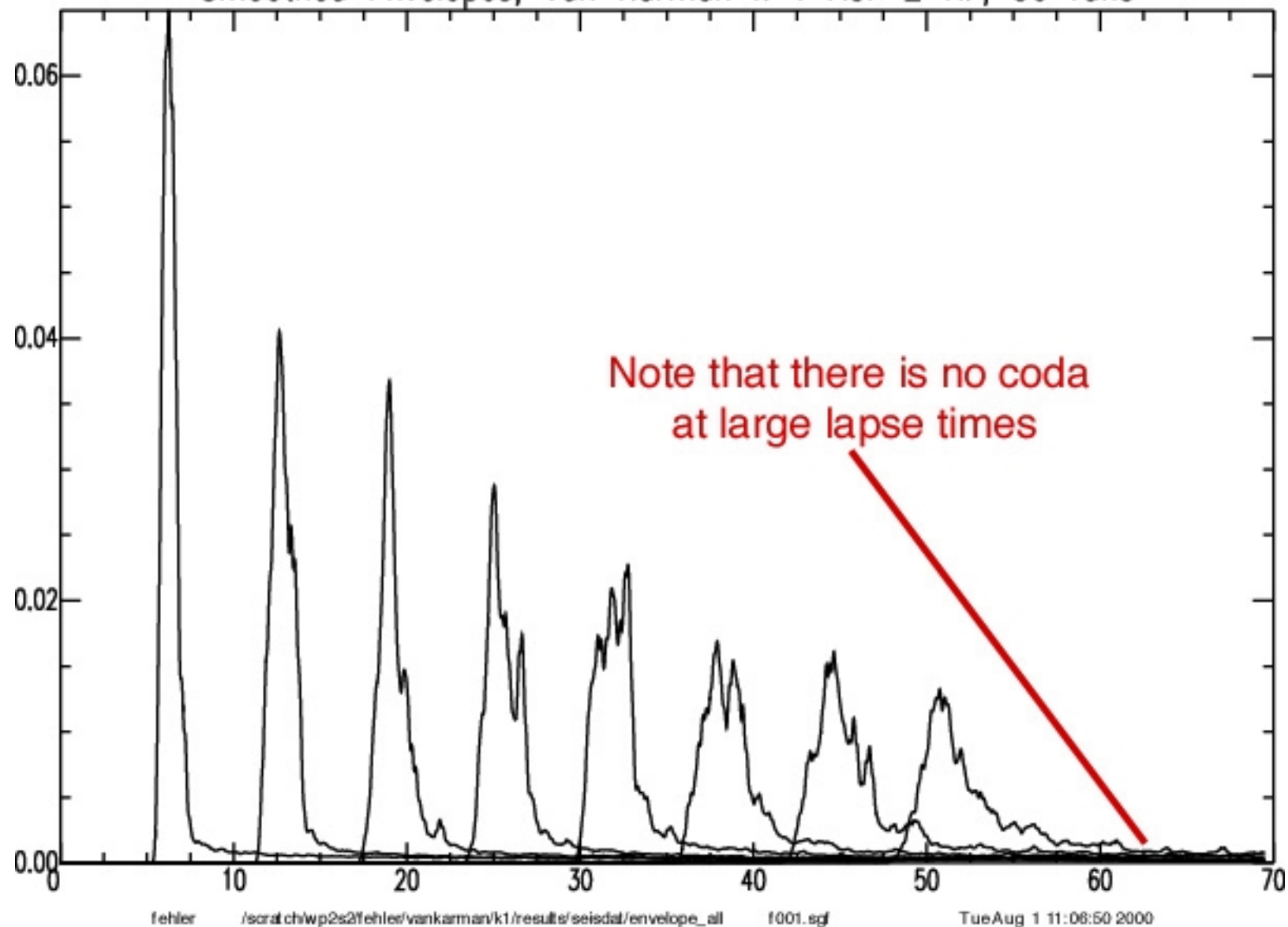
Gaussian ACF: $a = 5$ km, $\varepsilon = 0.05$, $V_0 = 4$ km/s



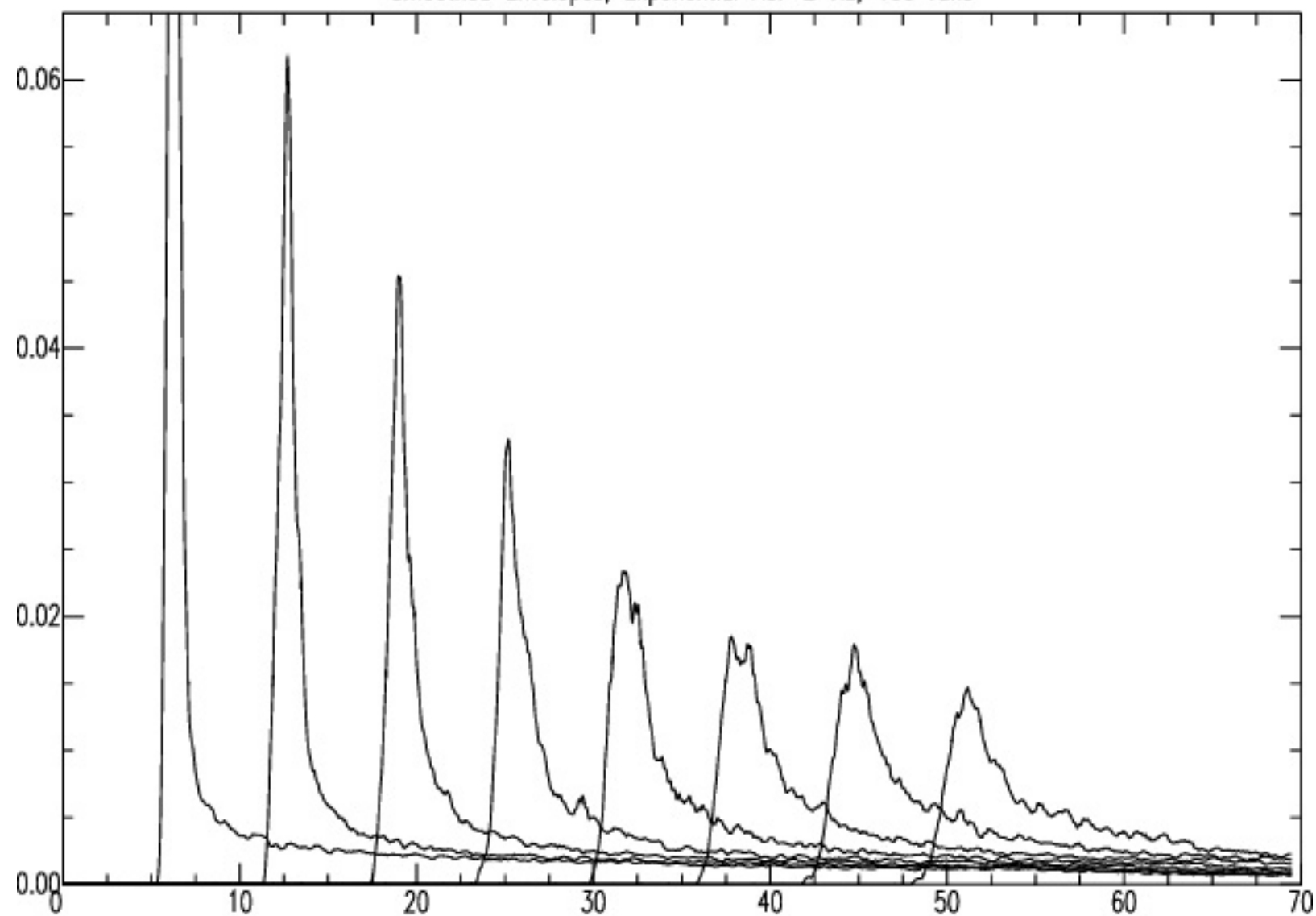
Comparisons with Different Correlation Functions

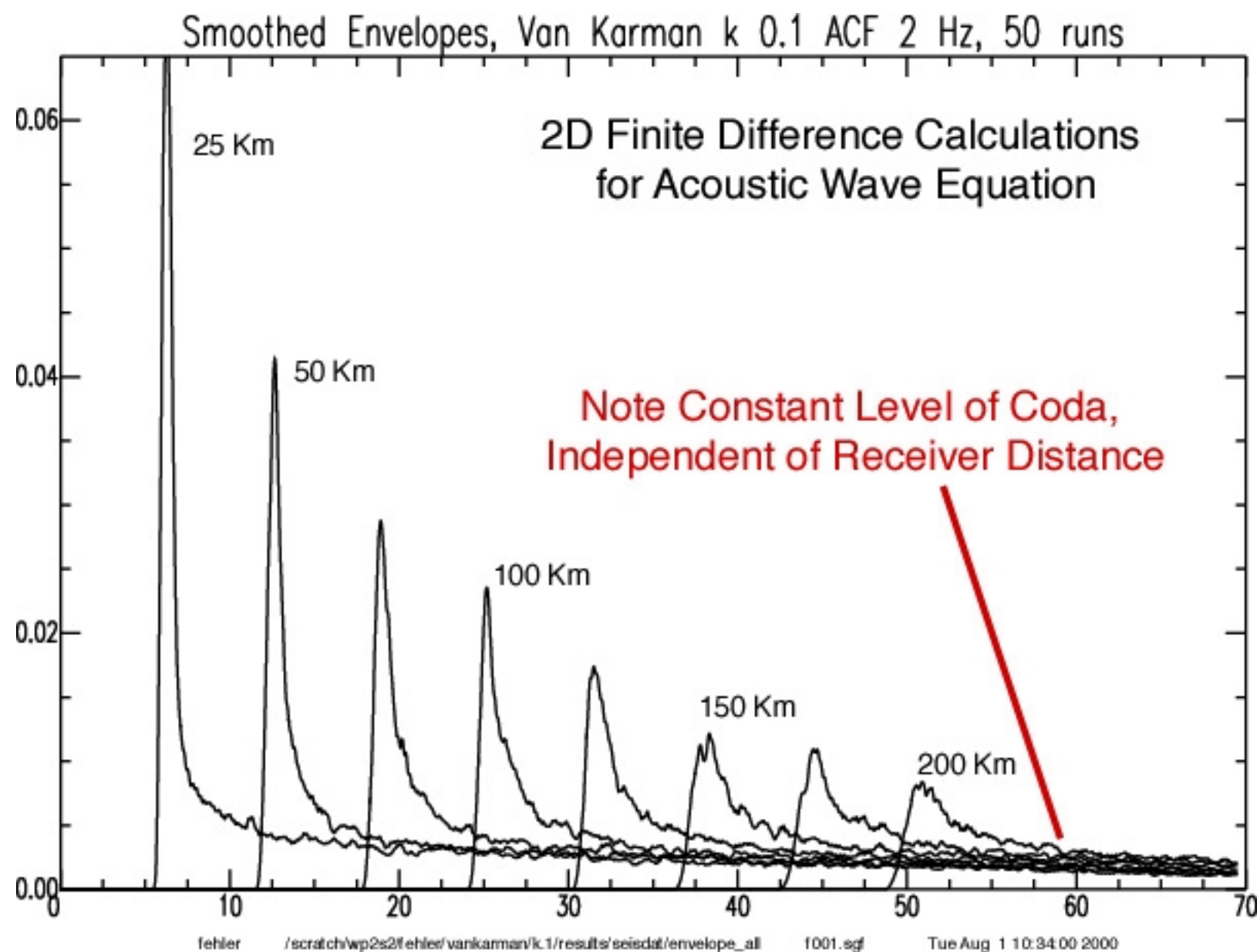


Smoothed Envelopes, Van Karman k 1 ACF 2 Hz, 50 runs

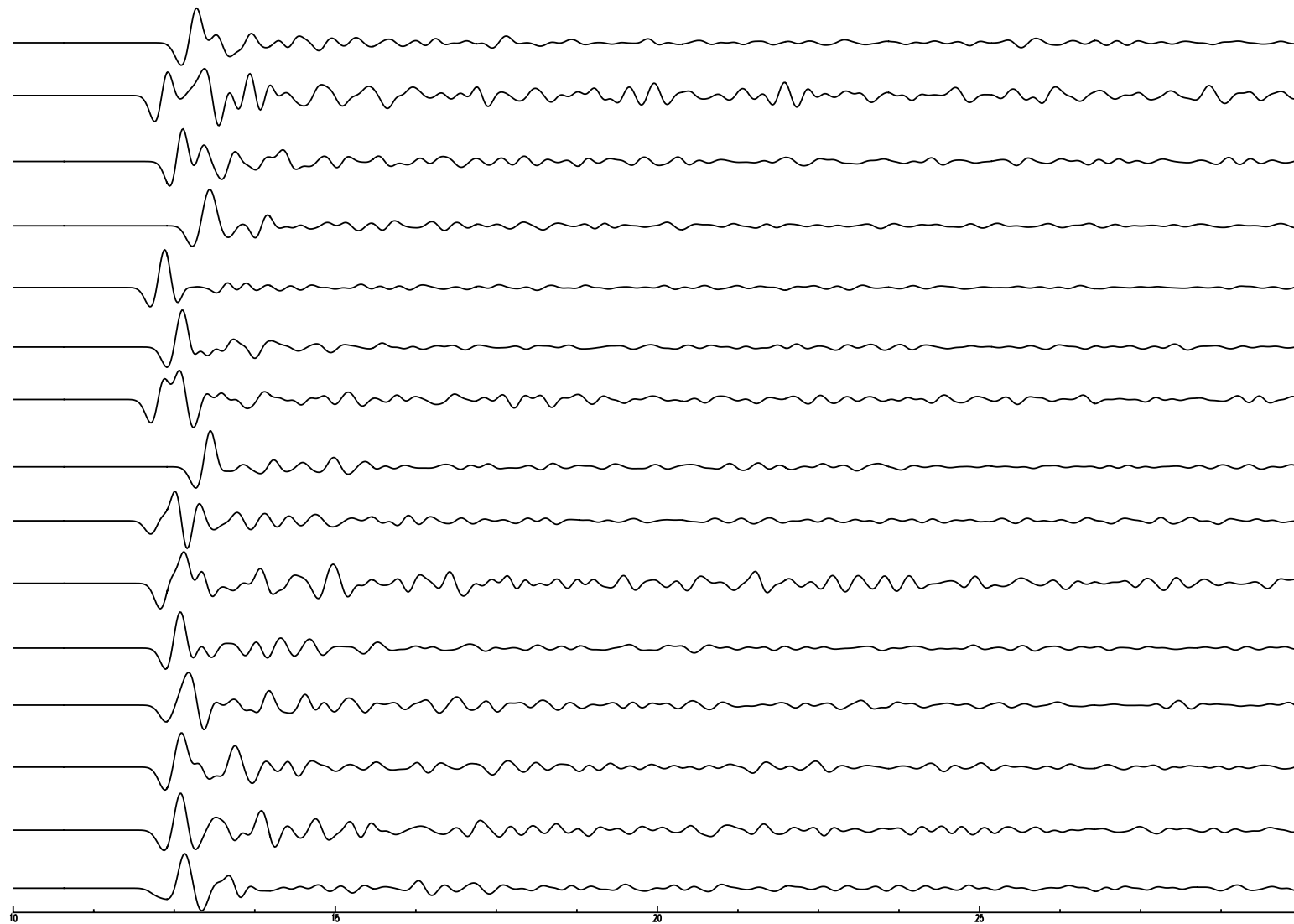


Smoothed Envelopes, Exponential ACF 2 Hz, 100 runs

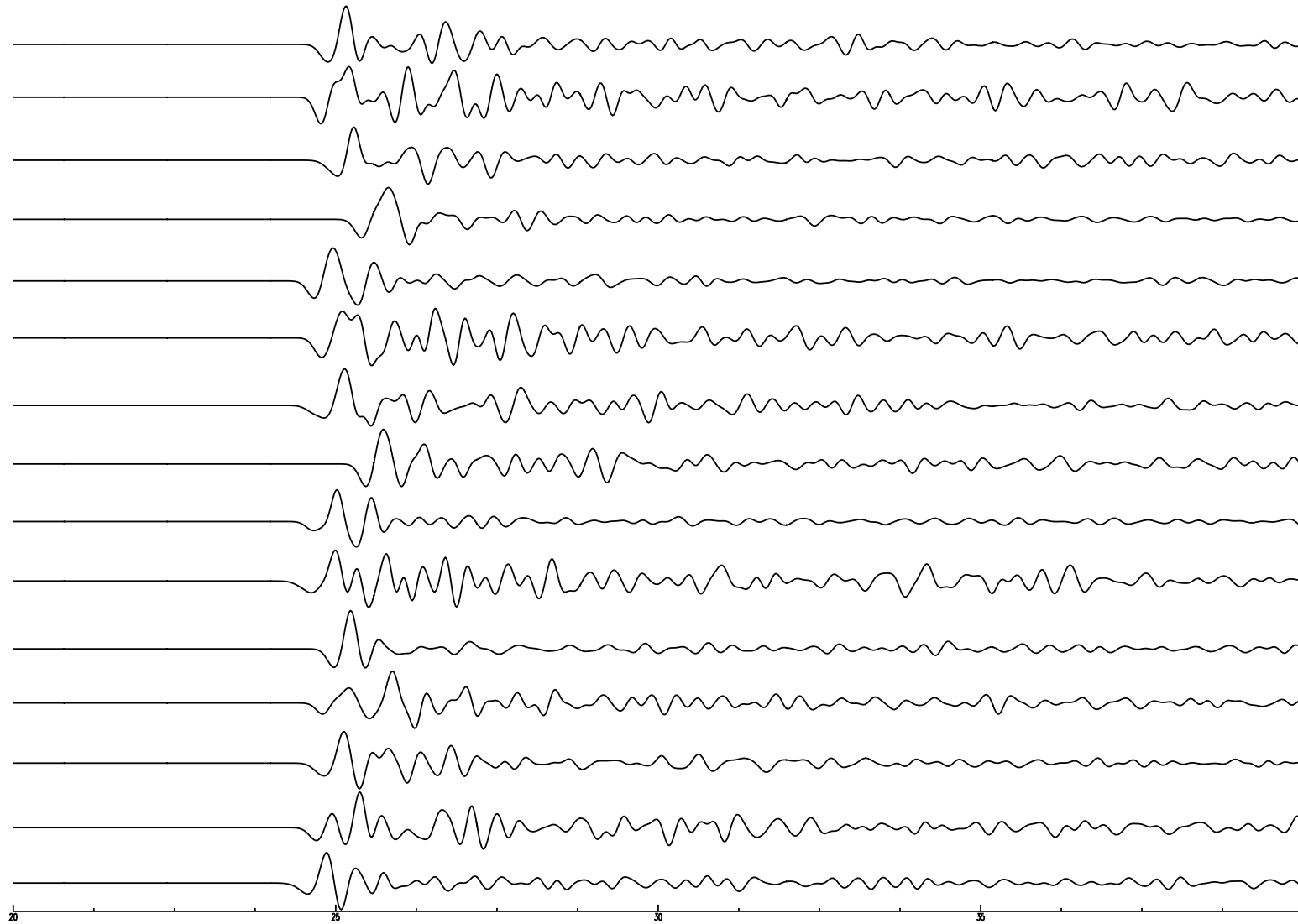




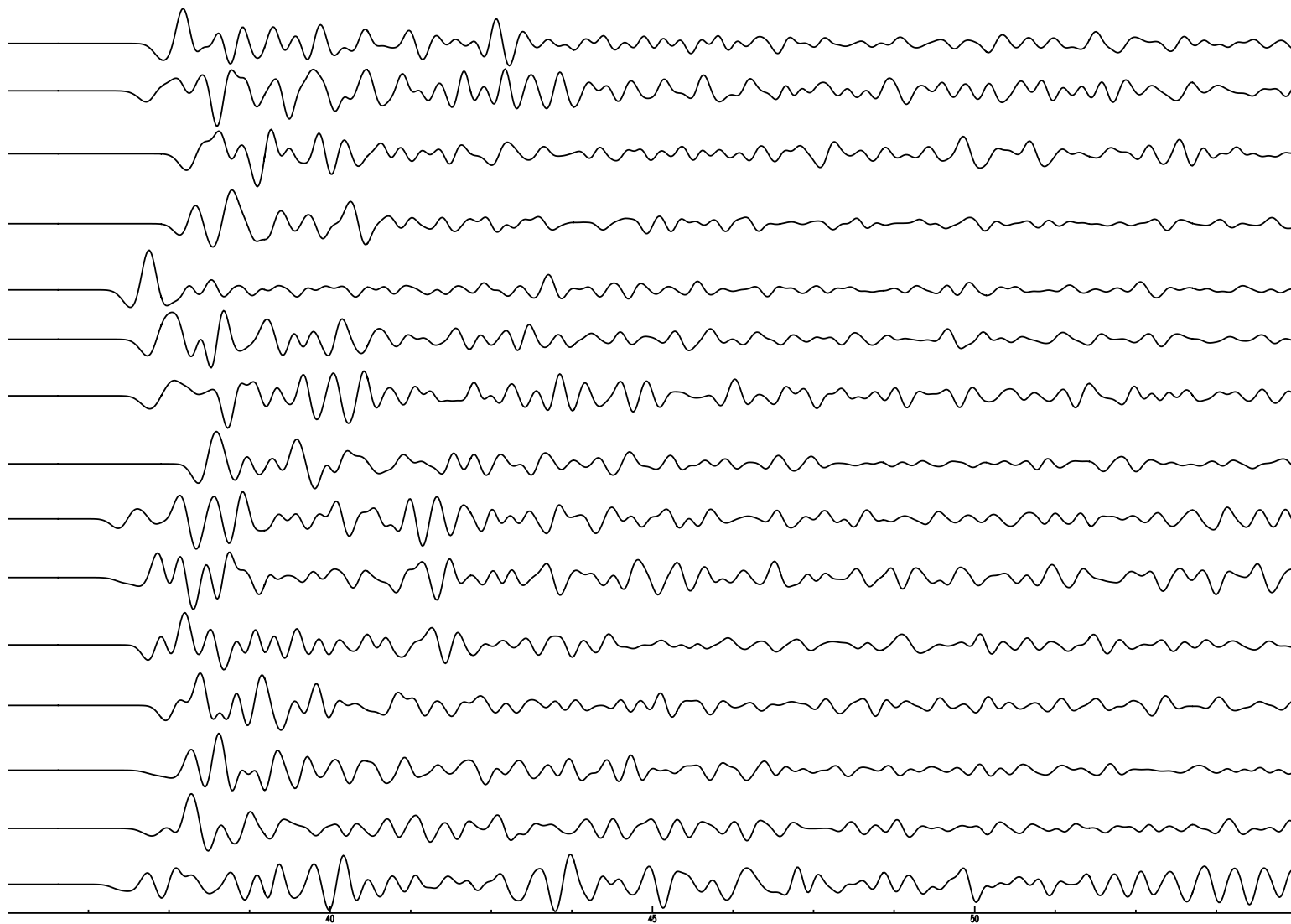
50 km 15 realizations, Van Karman k =.1 ACF, 2 Hz, Reduced Time



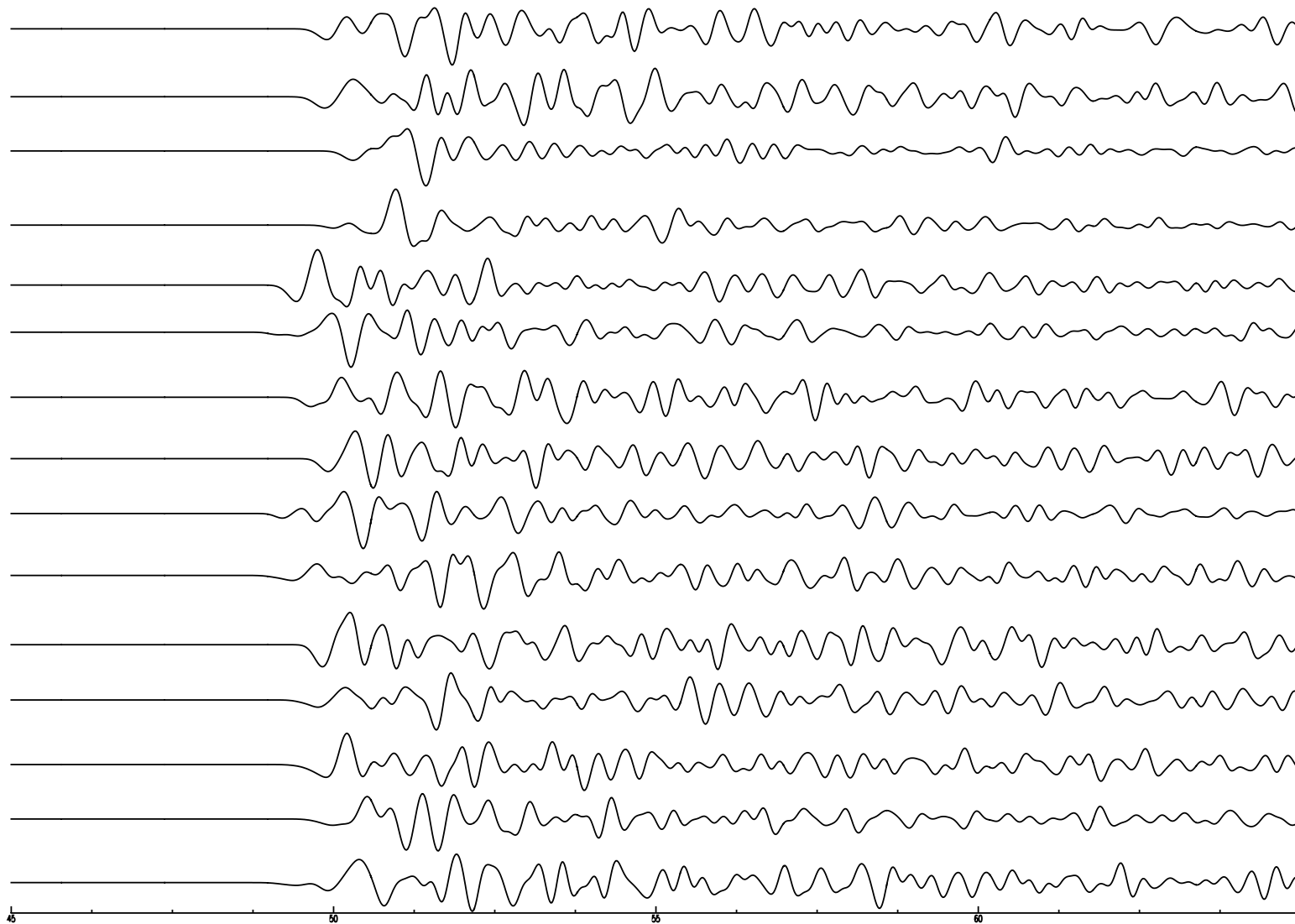
100 km 15 realizations, Van Karman $k = .1$ ACF, 2 Hz,



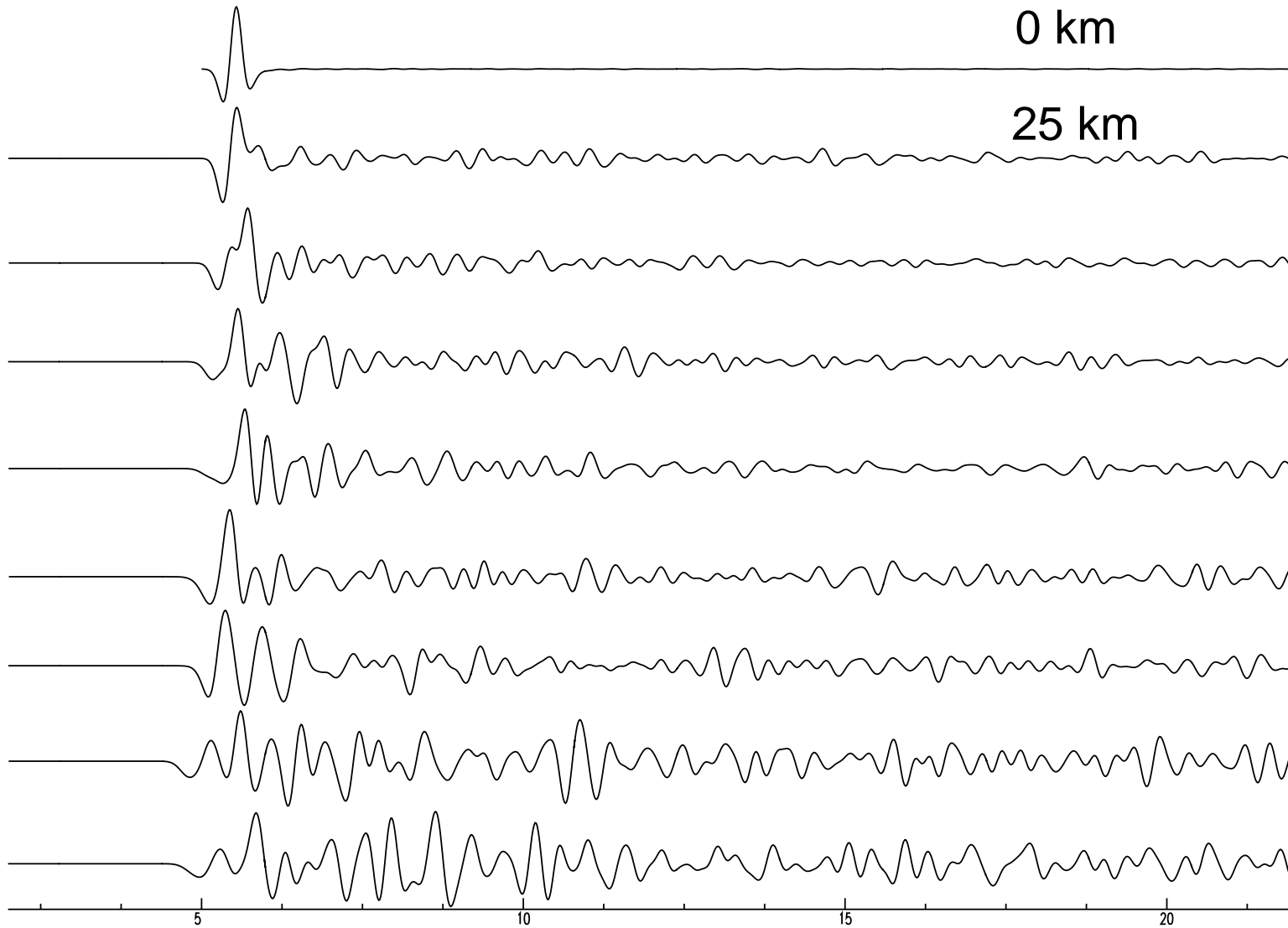
150 km 15 realizations, Van Karman $k = .1$ ACF, 2 Hz



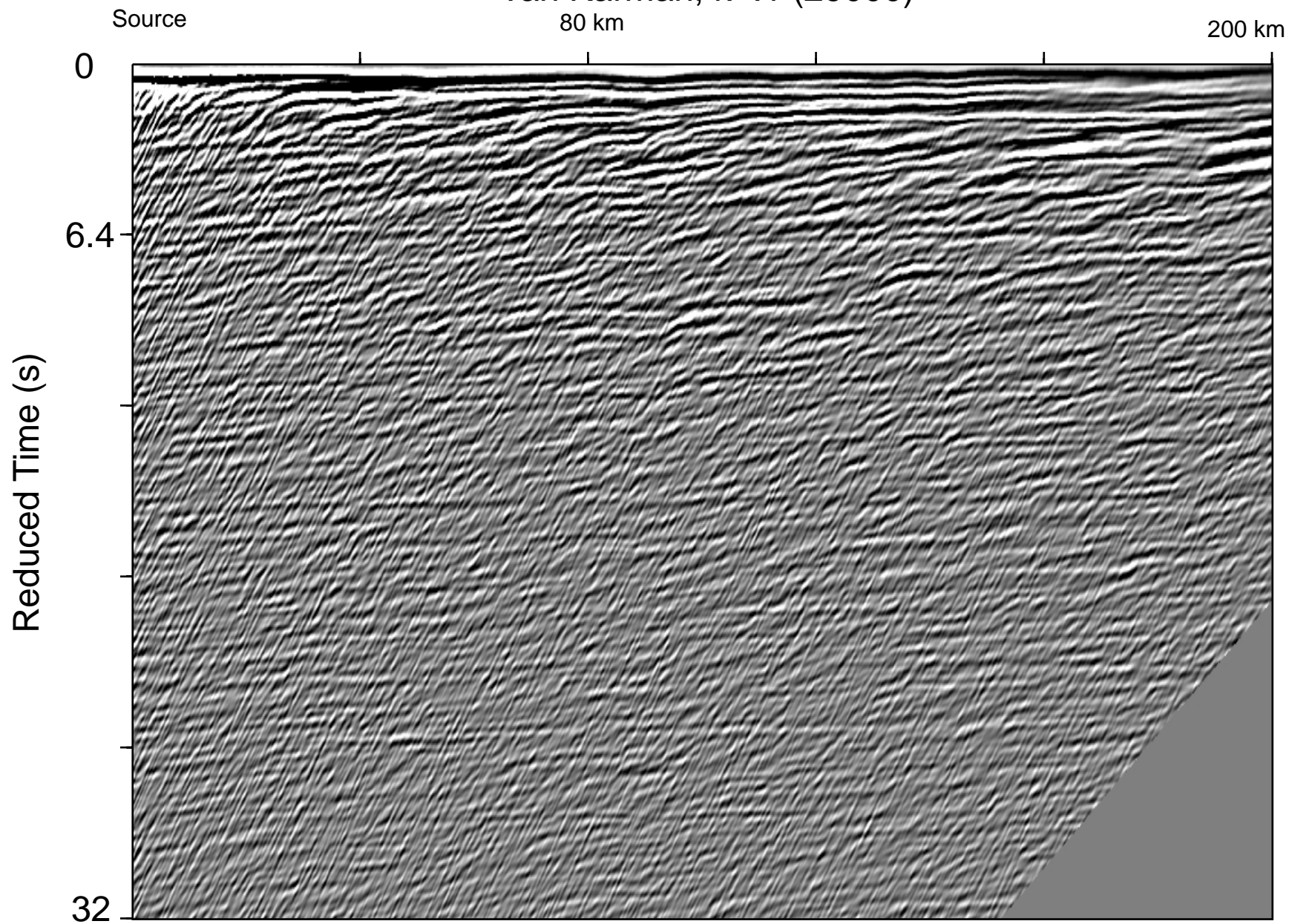
200 km 15 realizations, Van Karman $k = .1$ ACF, 2 Hz



Record Section k = 0.1

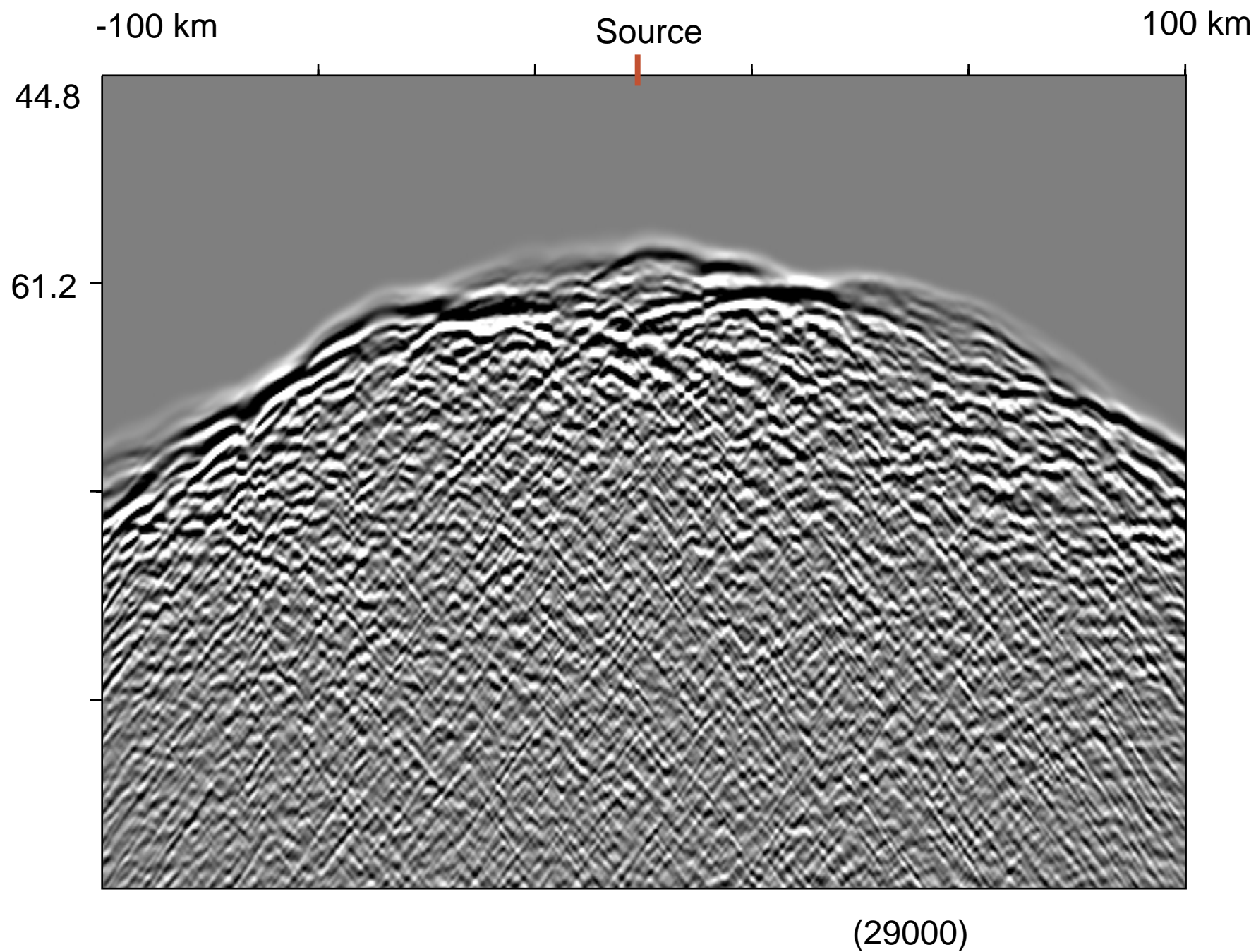


Van Karman, $k=.1$ (29000)

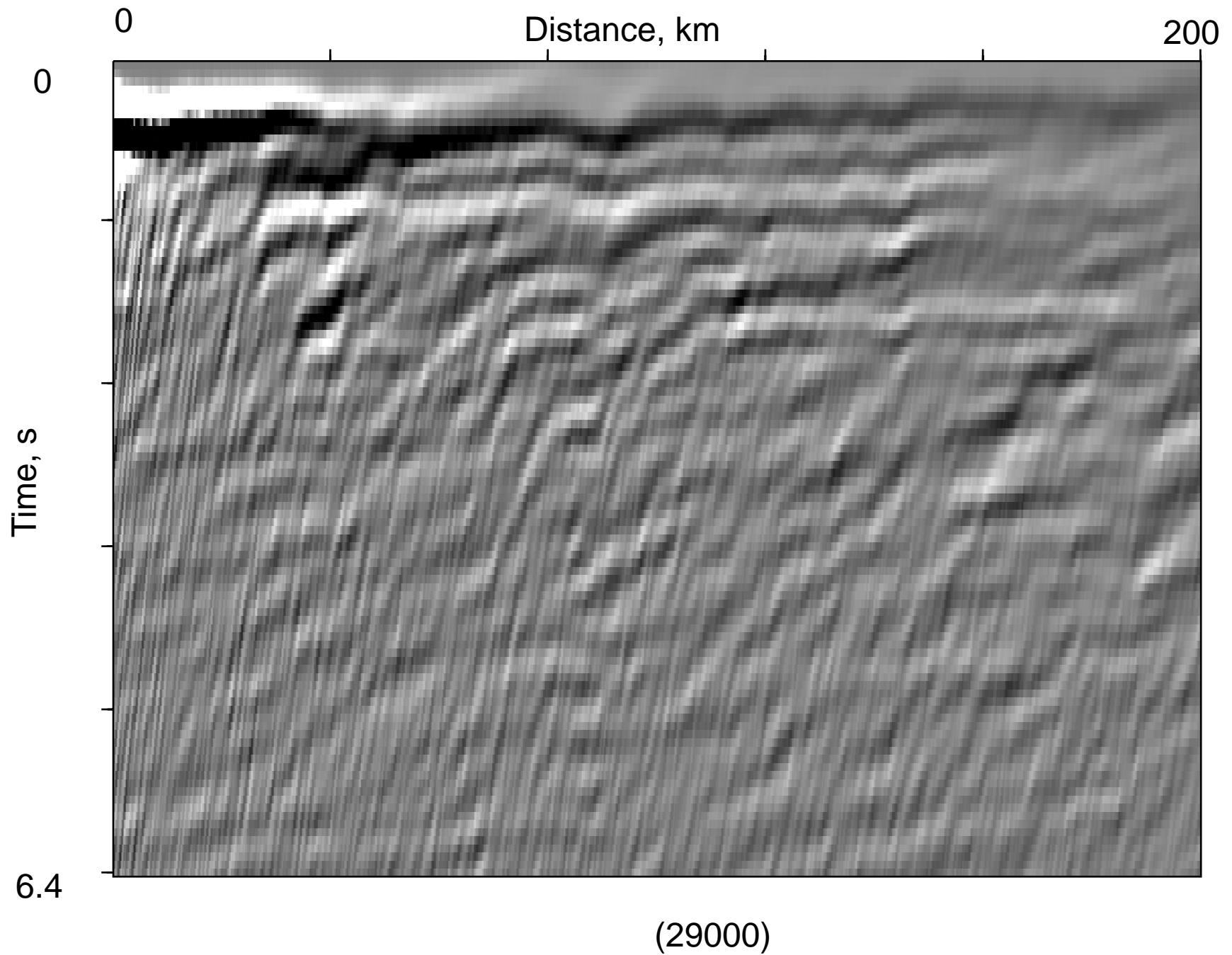


Record Section along propagation direction

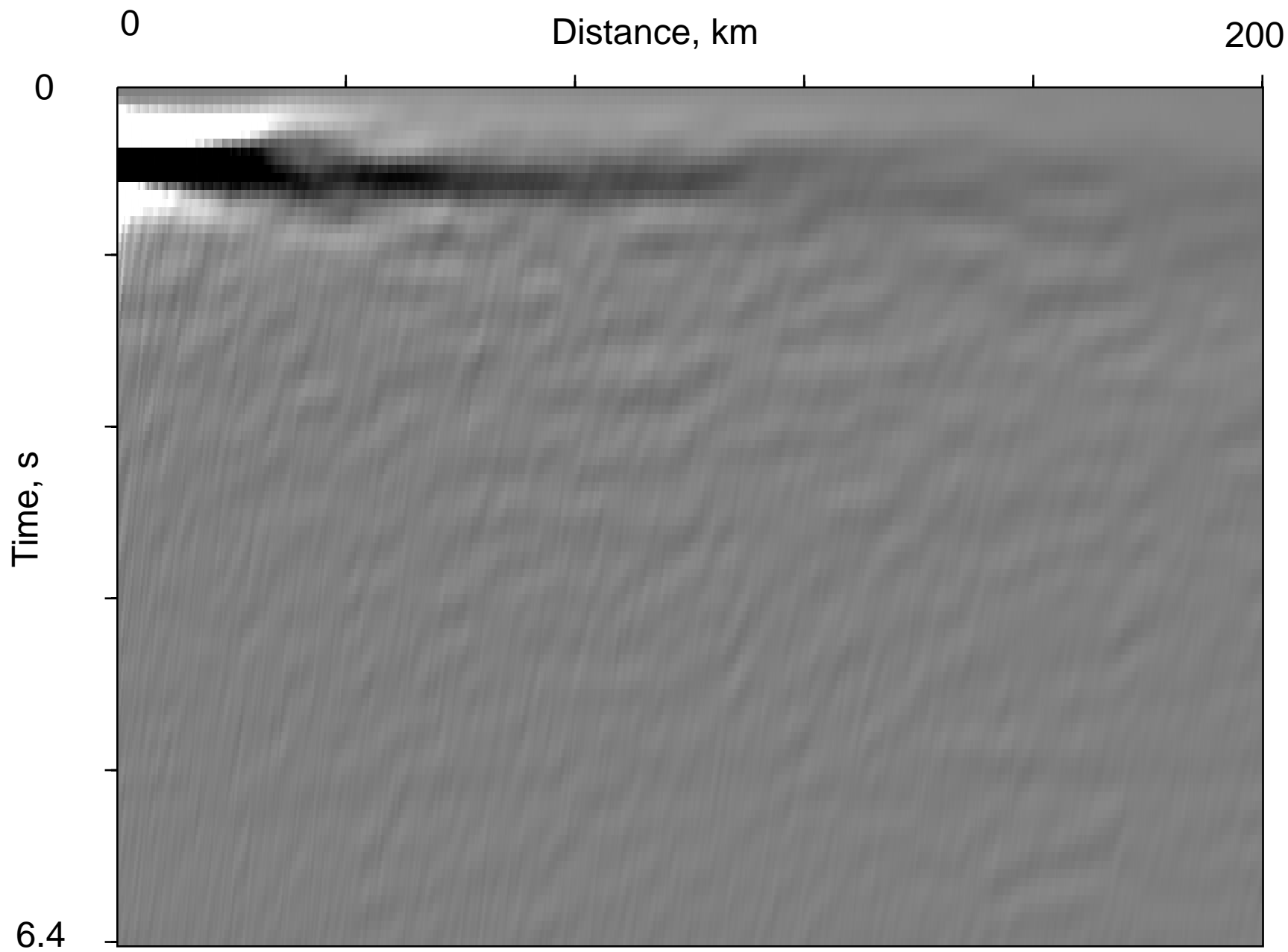
Van Karman, $k=0.1$, Receivers at $iz=200$ km from source



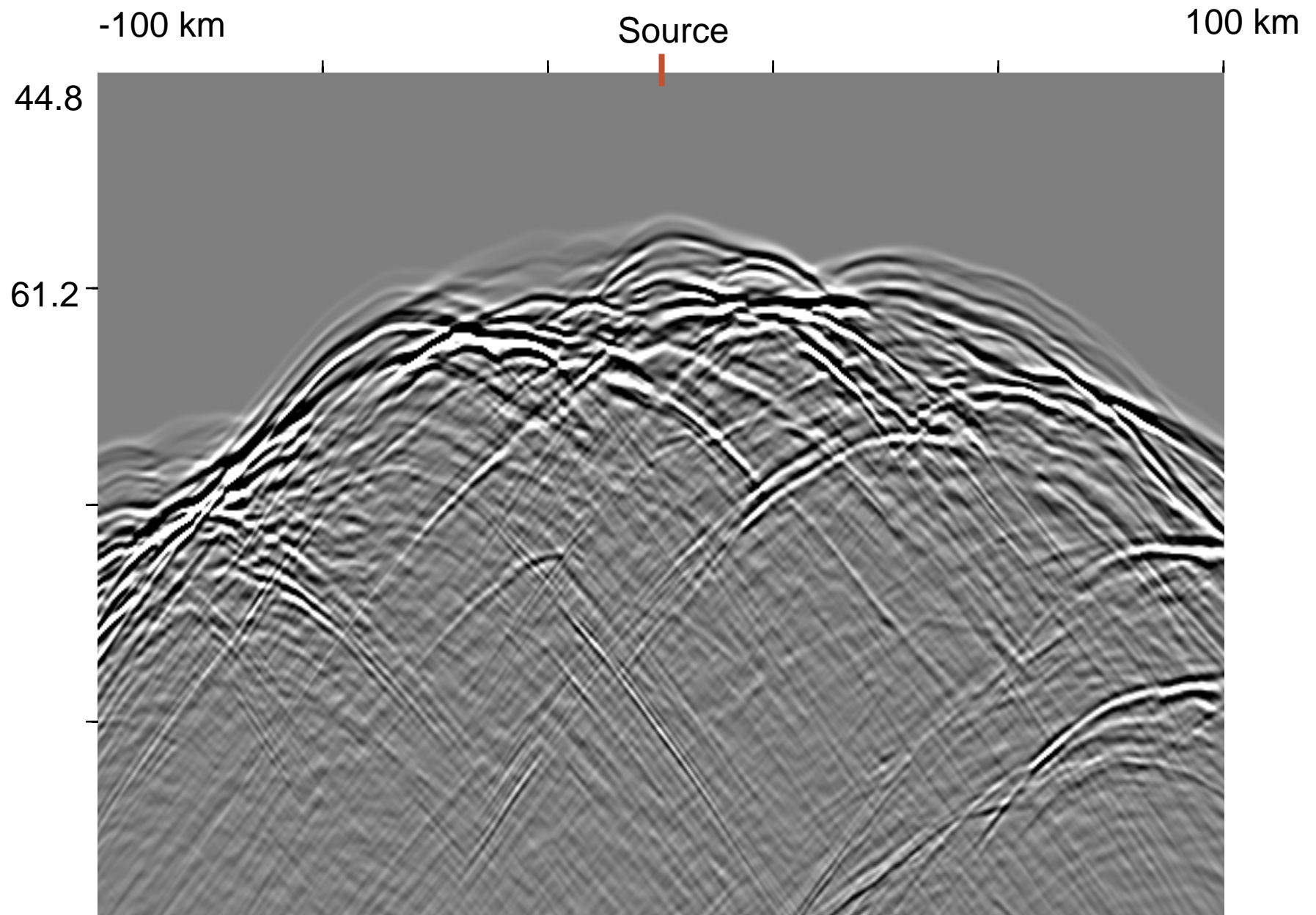
Single relization, $k=0.1$ van Karman



Stack Traces from 50 relizations, $k=0.1$ van Karman

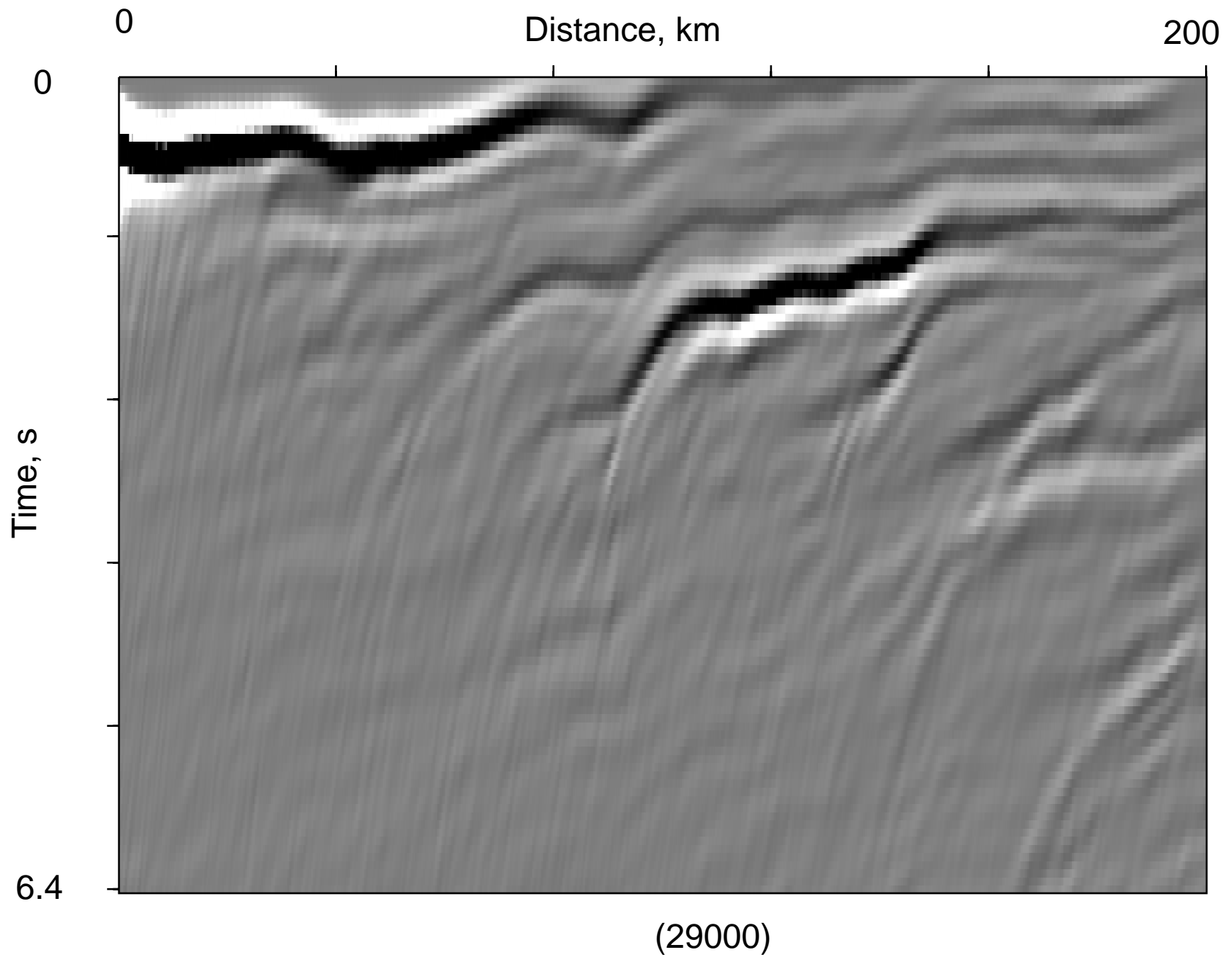


Van Karman, $k=1.$, Receivers at $iz=200$ km from source

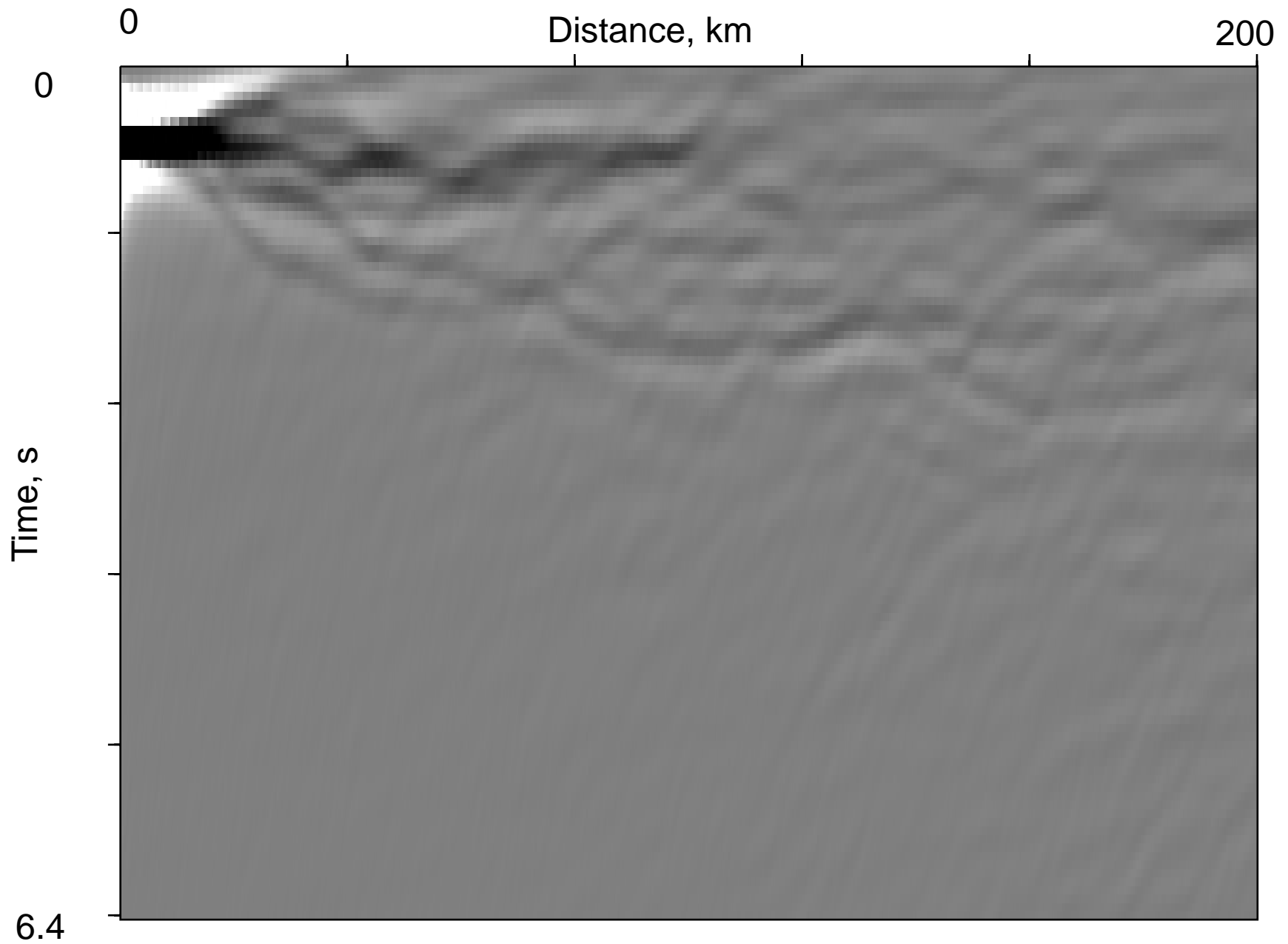


(29000)

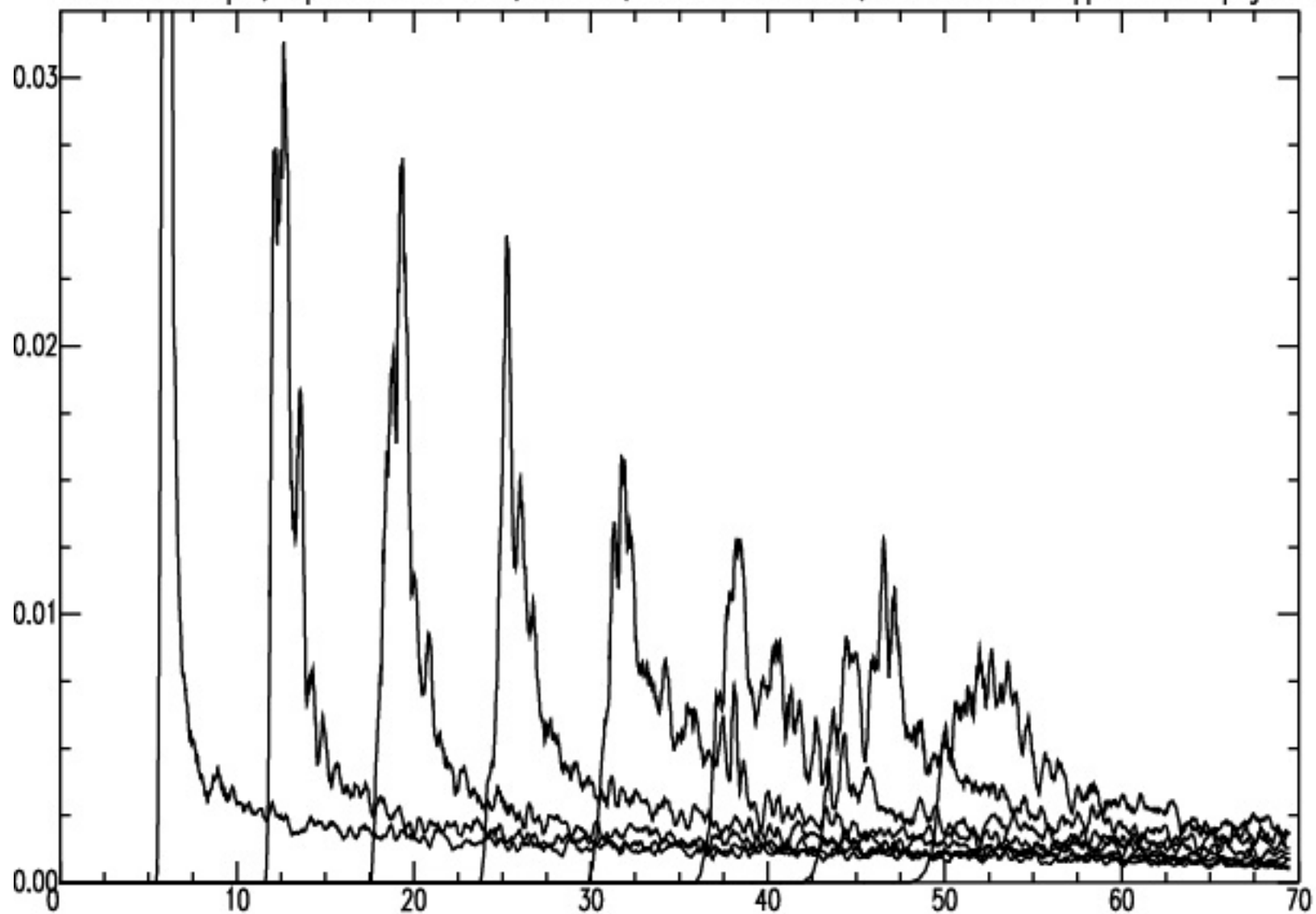
Single relization, k=1. van Karman



Stack Traces from 40 relizations, k=1. van Karman



Smoothed Envelopes, Exponential ACF 4 Hz, 10 Runs, times 10 to correct, 4 Hz correction applied as on page 33



Conclusions

- At close range, waveforms vary mostly due to traveltime fluctuation; at large range, scattering becomes important
- Good Agreement of Markov and Extended Local Rytov Results
- Finite Difference Results are a little different than Approximate Results
 - Importance of wide angle scattering

Conclusions

- Backscattering from near-source and near-receiver regions is important

